

Through The Looking Glass:

Faraday Conversion
in
Turbulent Blazar Jets



Nicholas MacDonald

Boston University

Talk Outline

A brief introduction to Faraday Conversion.

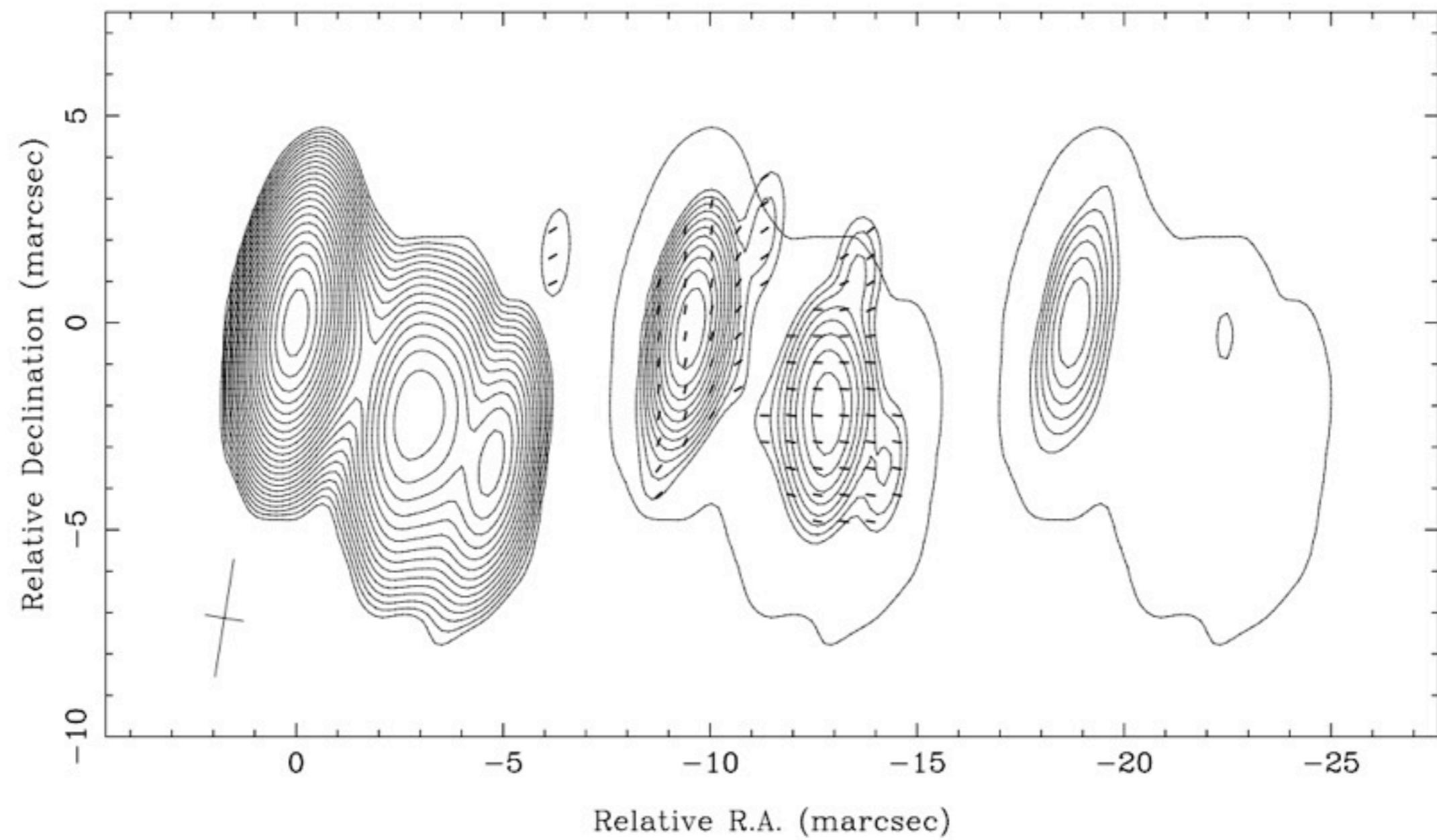
Can circular polarization be produced within a turbulent jet?

Malaga, Spain
June 2016



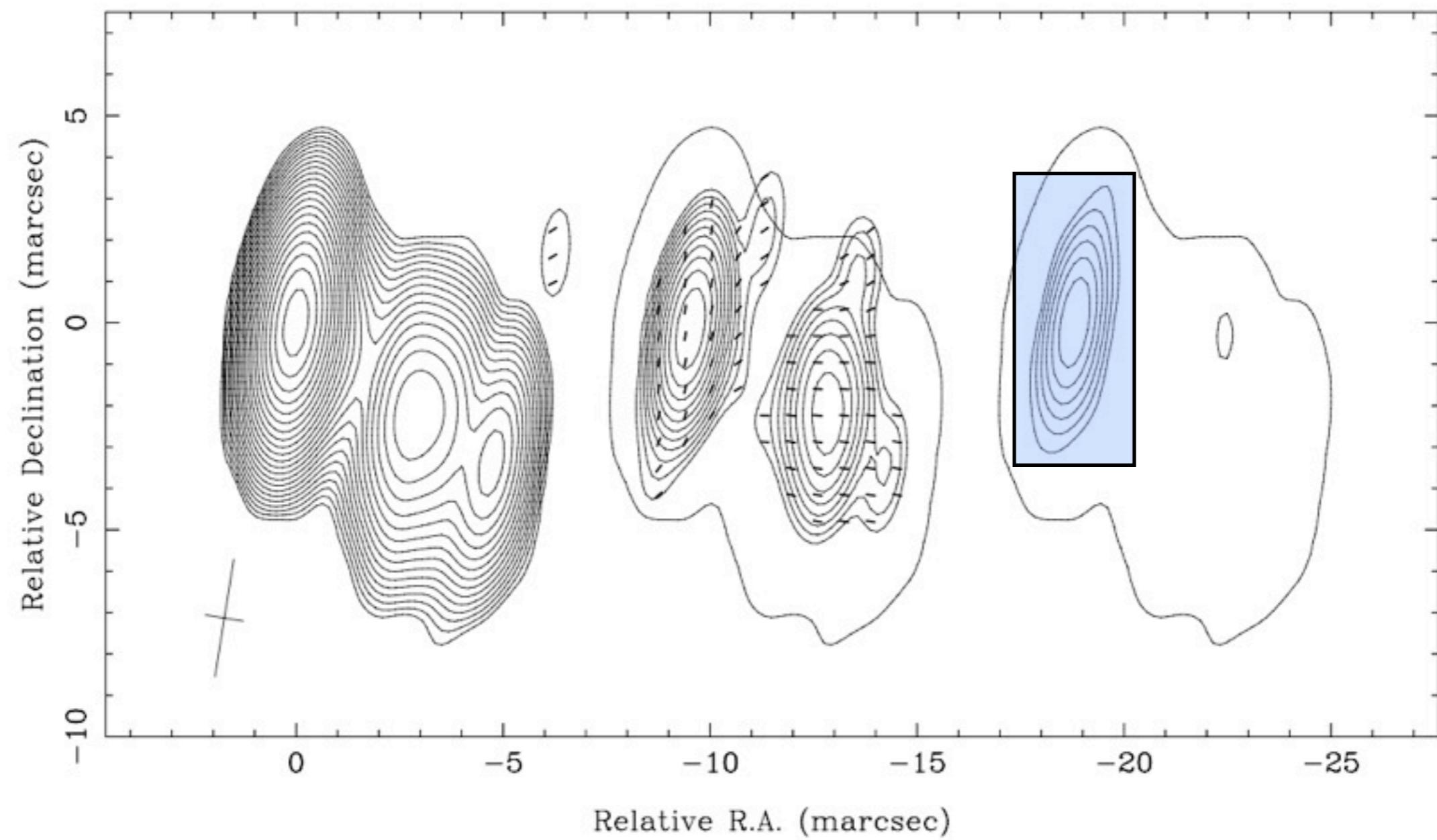
Birefringence

3C 279



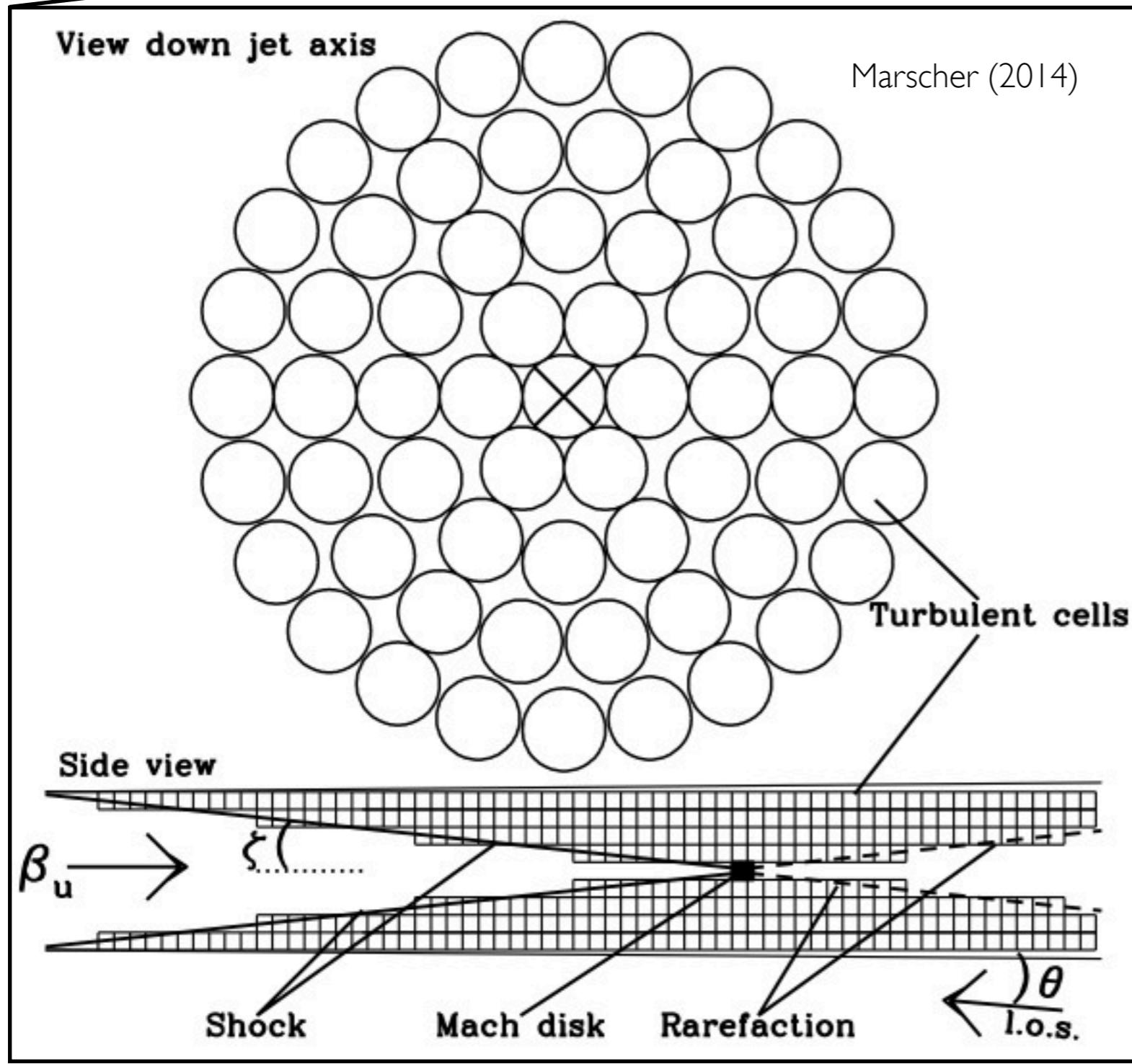
Homan et al. (2009)

3C 279

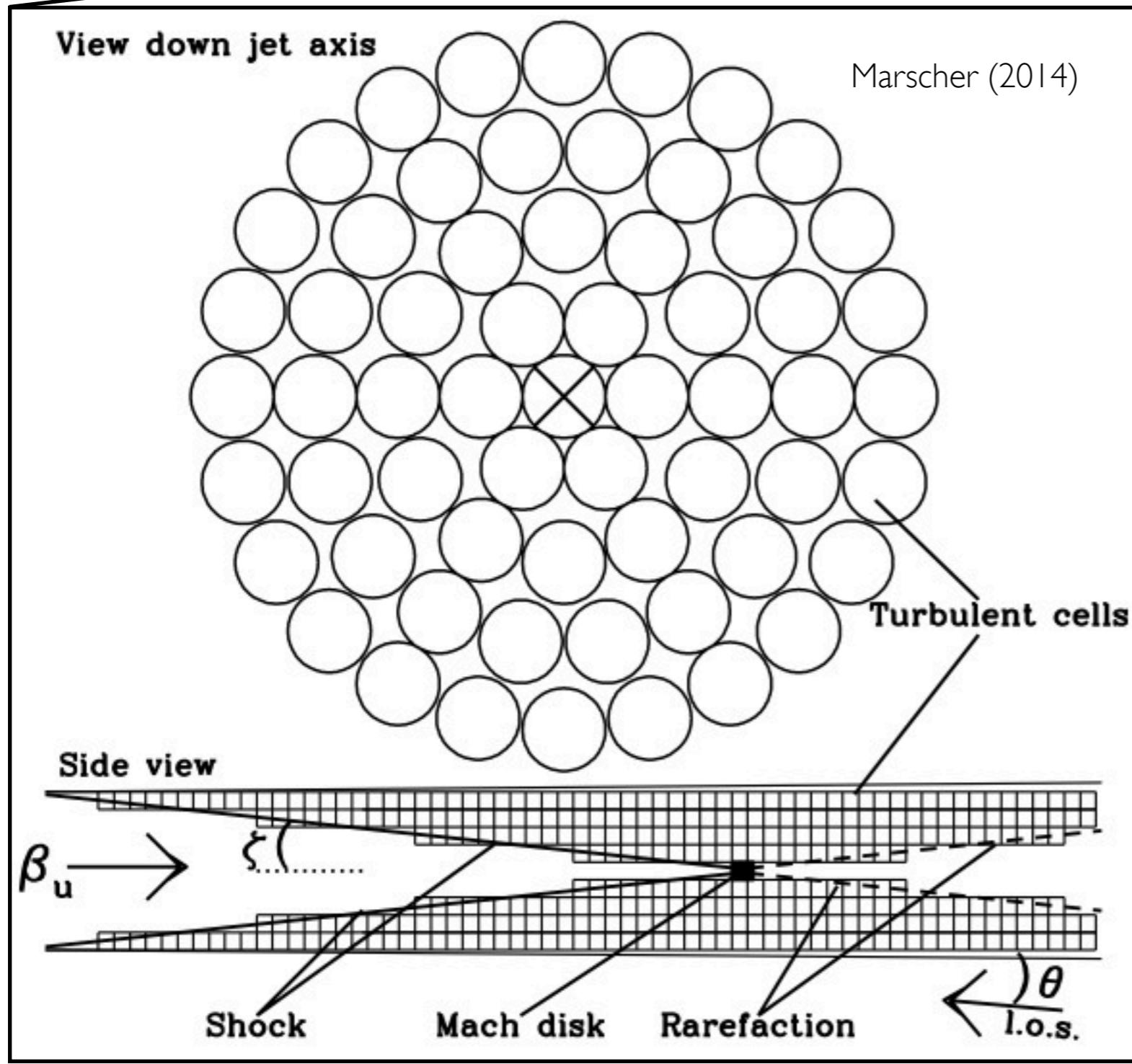


Homan et al. (2009)

The turbulent extreme multi-zone (TEMZ) code:



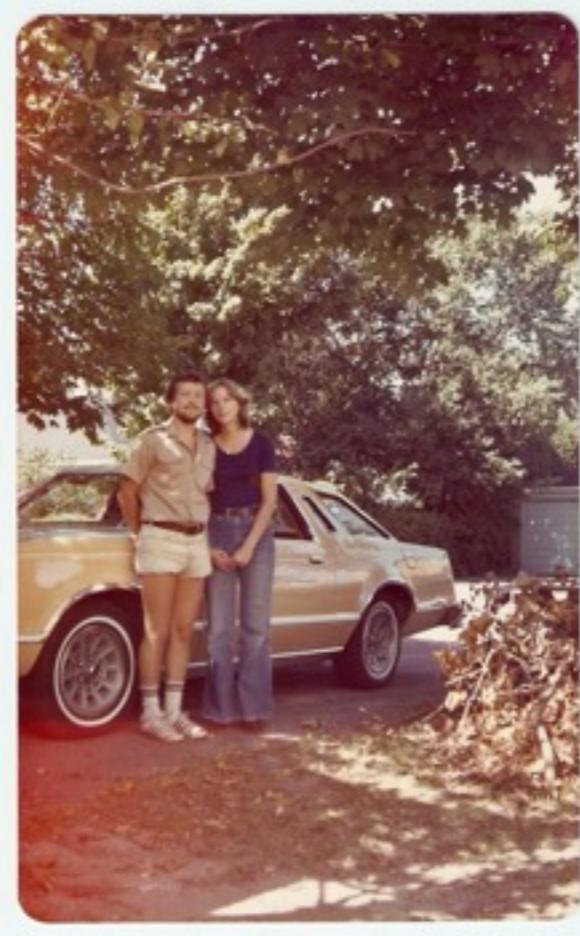
The turbulent extreme multi-zone (TEMZ) code:



?

1977

1977





1977



TRANSFER OF POLARIZED RADIATION IN SELF-ABSORBED SYNCHROTRON SOURCES.
I. RESULTS FOR A HOMOGENEOUS SOURCE

T. W. JONES

National Radio Astronomy Observatory*

AND

S. L. O'DELL

Department of Physics, University of California, San Diego

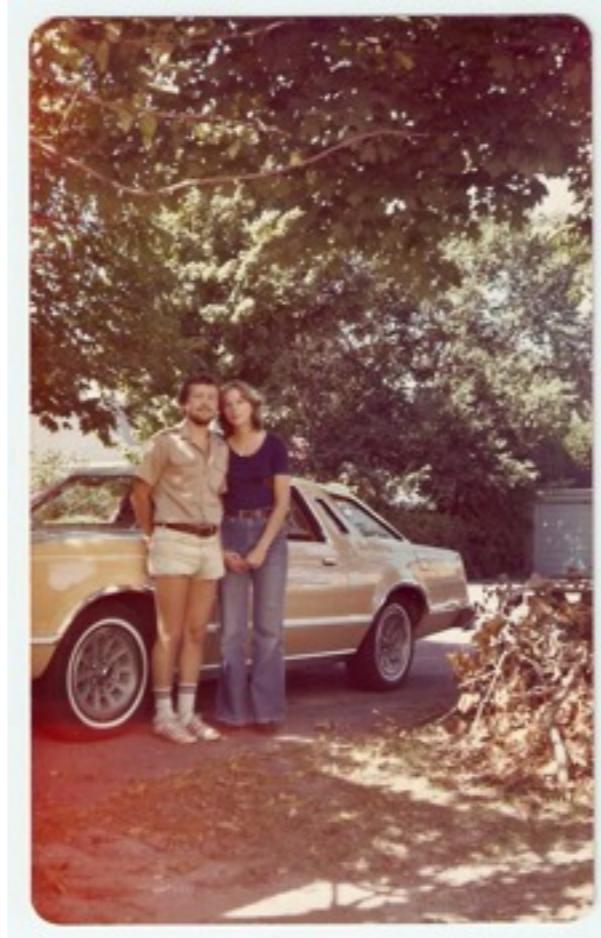
Received 1975 June 16; revised 1976 October 26

ABSTRACT

The solution to the equation of transfer of polarized radiation in a stationary, homogeneous, rarefied medium is applied to self-absorbed synchrotron sources. Relativistic electrons (independent of the presence of any cold plasma) can quite easily produce in such sources significant Faraday rotation and/or conversion of linear to circular polarization. Structural inhomogeneities do not obviate the importance of these phenomena in cosmic, compact nonthermal sources. Contrary to the calculation of Pacholczyk and Swihart, the circular polarization for a homogeneous source changes sign just below the self-absorption turnover as the source becomes opaque, even when polarization conversion dominates; however, for a physically realistic source, structural inhomogeneity may alter this behavior. The observational evidence bearing upon these effects is reviewed.

Subject headings: polarization — radiative transfer — radio sources: variable — synchrotron radiation

1977



TRANSFER OF POLARIZED RADIATION IN SELF-ABSORBED SYNCHROTRON SOURCES.
I. RESULTS FOR A HOMOGENEOUS SOURCE

T. W. JONES

National Radio Astronomy Observatory*

AND

S. L. O'DELL

Department of Physics, University of California, San Diego

Received 1975 June 16; revised 1976 October 26

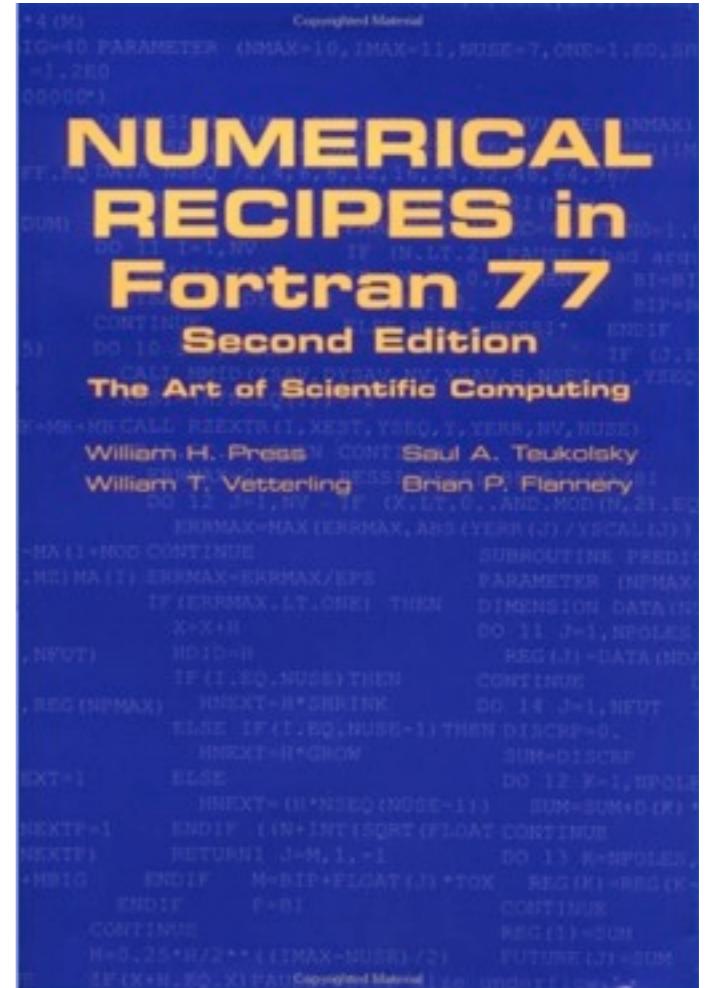
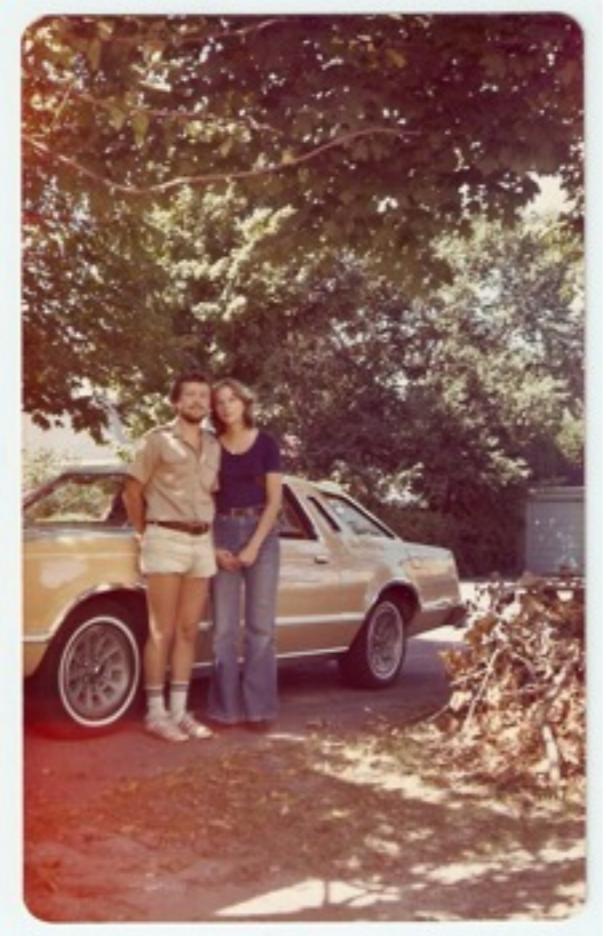
ABSTRACT

The solution to the equation of transfer of polarized radiation in a stationary, homogeneous, rarefied medium is applied to self-absorbed synchrotron sources. Relativistic electrons (independent of the presence of any cold plasma) can quite easily produce in such sources significant Faraday rotation and/or conversion of linear to circular polarization. Structural inhomogeneities do not obviate the importance of these phenomena in cosmic, compact nonthermal sources. Contrary to the calculation of Pacholczyk and Swihart, the circular polarization for a homogeneous source changes sign just below the self-absorption turnover as the source becomes opaque, even when polarization conversion dominates; however, for a physically realistic source, structural inhomogeneity may alter this behavior. The observational evidence bearing upon these effects is reviewed.

Subject headings: polarization — radiative transfer — radio sources: variable — synchrotron radiation



1977



TRANSFER OF POLARIZED RADIATION IN SELF-ABSORBED SYNCHROTRON SOURCES.
I. RESULTS FOR A HOMOGENEOUS SOURCE

T. W. JONES

National Radio Astronomy Observatory*

AND

S. L. O'DELL

Department of Physics, University of California, San Diego

Received 1975 June 16; revised 1976 October 26

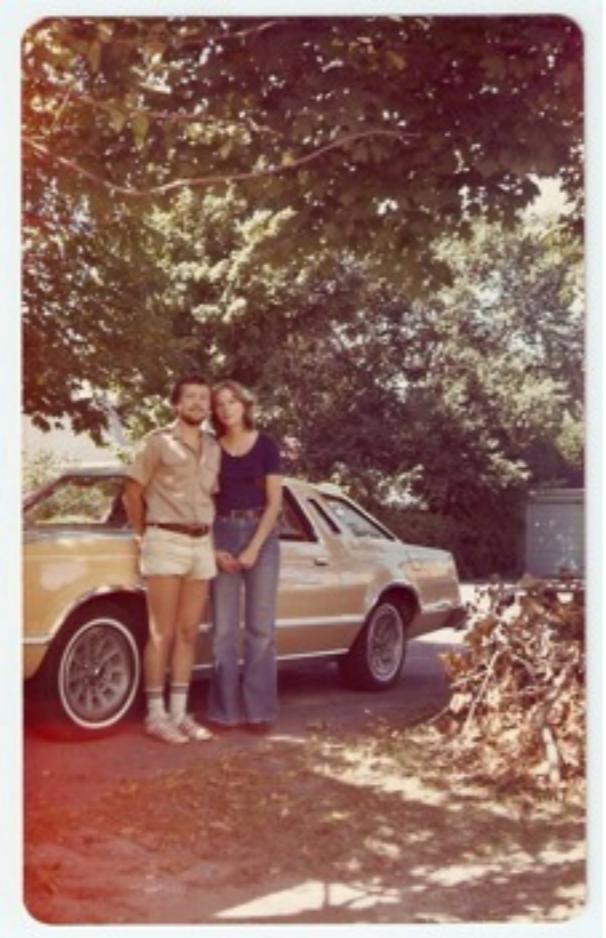
ABSTRACT

The solution to the equation of transfer of polarized radiation in a stationary, homogeneous, rarefied medium is applied to self-absorbed synchrotron sources. Relativistic electrons (independent of the presence of any cold plasma) can quite easily produce in such sources significant Faraday rotation and/or conversion of linear to circular polarization. Structural inhomogeneities do not obviate the importance of these phenomena in cosmic, compact nonthermal sources. Contrary to the calculation of Pacholczyk and Swihart, the circular polarization for a homogeneous source changes sign just below the self-absorption turnover as the source becomes opaque, even when polarization conversion dominates; however, for a physically realistic source, structural inhomogeneity may alter this behavior. The observational evidence bearing upon these effects is reviewed.

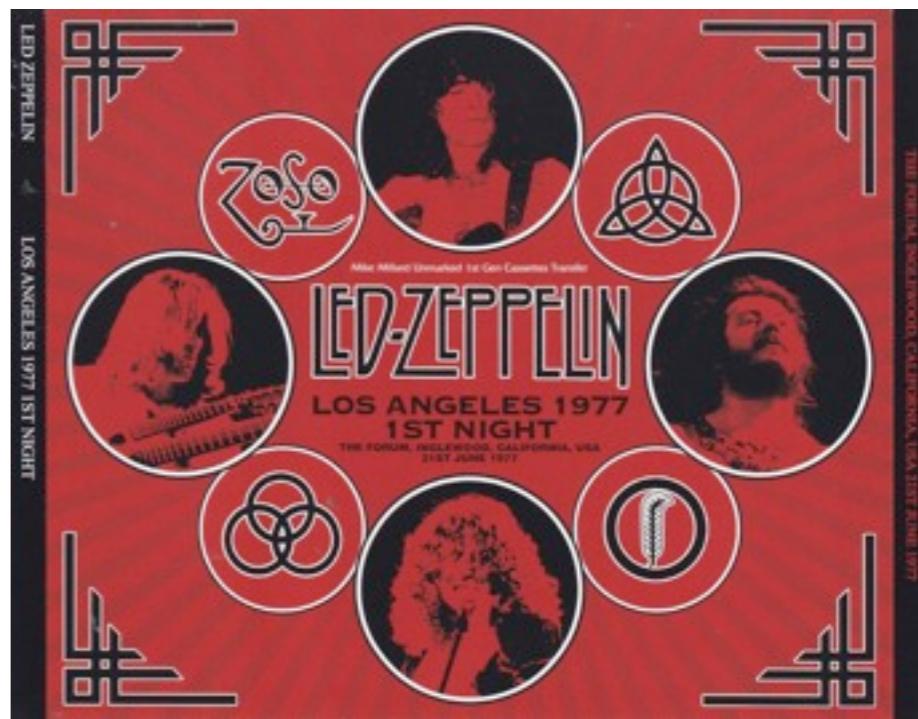
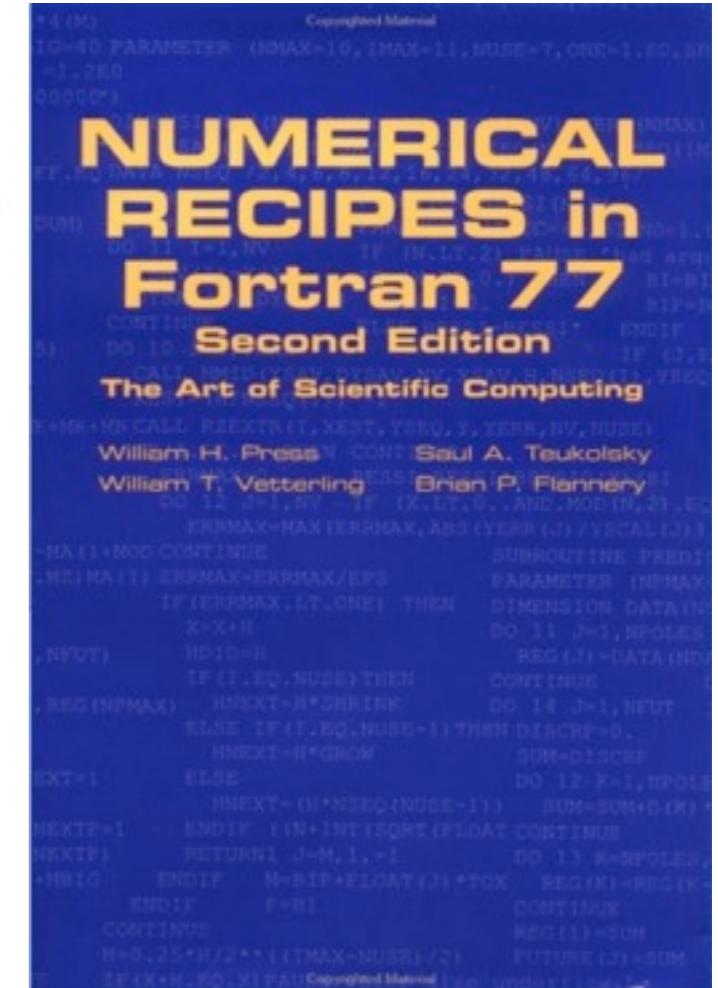
Subject headings: polarization — radiative transfer — radio sources: variable — synchrotron radiation



1977



6



Solving the Full Stokes Equations of Radiative Transfer

$$\begin{array}{c|c|c|c|c} \left(\frac{d}{dl} + \kappa_I \right) & \kappa_Q & \kappa_U & \kappa_V & I_\nu \\ \kappa_Q & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_V & -\kappa^*_U & Q_\nu \\ \kappa_U & -\kappa^*_V & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_Q & U_\nu \\ \kappa_V & \kappa^*_U & -\kappa^*_Q & \left(\frac{d}{dl} + \kappa_I \right) & V_\nu \end{array} = \begin{array}{c} \eta_\nu^I \\ \eta_\nu^Q \\ \eta_\nu^U \\ \eta_\nu^V \end{array}$$

Jones & O'Dell (1977)

Stokes Parameters

$$\begin{pmatrix} \frac{d}{dl} + \kappa_I \\ \kappa_Q \\ \kappa_U \\ \kappa_V \end{pmatrix} = \begin{pmatrix} \frac{d}{dl} + \kappa_I \\ \kappa_Q \\ -\kappa^*_V \\ \kappa^*_U \end{pmatrix}$$

$$\begin{pmatrix} \kappa_U \\ \kappa^*_V \\ -\kappa^*_Q \\ \frac{d}{dl} + \kappa_I \end{pmatrix} = \begin{pmatrix} \kappa_V \\ -\kappa^*_U \\ \kappa^*_Q \\ \frac{d}{dl} + \kappa_I \end{pmatrix}$$

$$= \begin{vmatrix} I_\nu \\ Q_\nu \\ U_\nu \\ V_\nu \end{vmatrix} = \begin{vmatrix} \eta_\nu^I \\ \eta_\nu^Q \\ \eta_\nu^U \\ \eta_\nu^V \end{vmatrix}$$

Jones & O'Dell (1977)

Emission Coefficients

$$\left(\frac{d}{dl} + \kappa_I \right)$$

$$\kappa_Q$$

$$\kappa_U$$

$$\kappa_V$$

$$\kappa_Q \quad \left(\frac{d}{dl} + \kappa_I \right)$$

$$-\kappa_V^*$$

$$\kappa_U^*$$

$$\kappa_U \quad \kappa_V^* \quad \left(\frac{d}{dl} + \kappa_I \right)$$

$$-\kappa_Q^*$$

$$\kappa_V \quad -\kappa_U^* \quad \kappa_Q^* \quad \left(\frac{d}{dl} + \kappa_I \right)$$

$$I_\nu$$

$$Q_\nu$$

$$U_\nu$$

$$V_\nu$$

$$\eta_\nu^I$$

$$\eta_\nu^Q$$

$$\eta_\nu^U$$

$$\eta_\nu^V$$

Jones & O'Dell (1977)

Absorption Coefficients

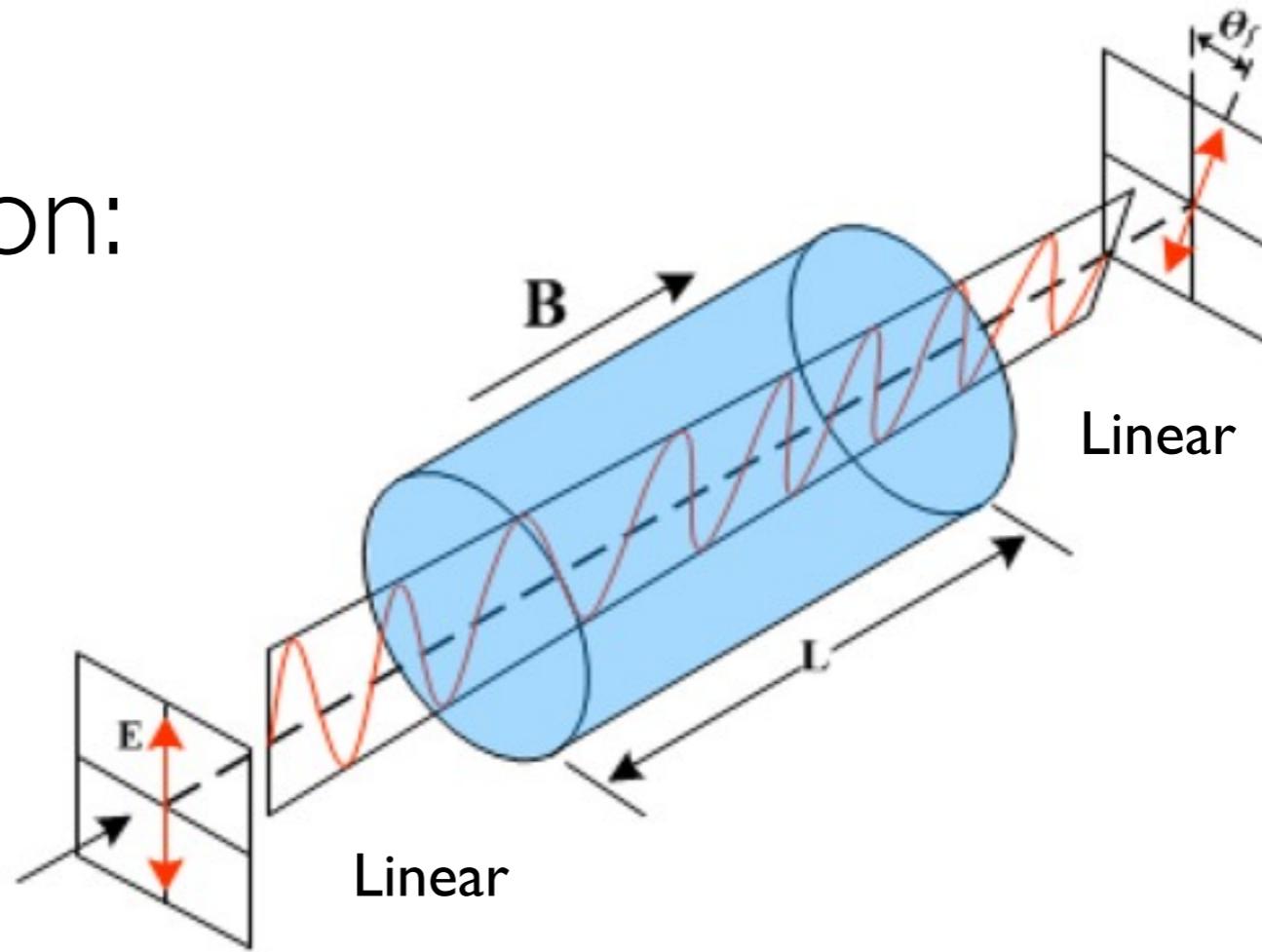
$$\begin{pmatrix}
 \frac{d}{dl} + \boxed{\kappa_I} \\
 \kappa_Q & \kappa_U & \kappa_V \\
 -\kappa^*_V & \frac{d}{dl} + \boxed{\kappa_I} & -\kappa^*_U \\
 \kappa^*_U & -\kappa^*_Q & \frac{d}{dl} + \boxed{\kappa_I}
 \end{pmatrix}
 \begin{vmatrix}
 I_\nu \\
 Q_\nu \\
 U_\nu \\
 V_\nu
 \end{vmatrix}
 = \begin{vmatrix}
 \eta_\nu^I \\
 \eta_\nu^Q \\
 \eta_\nu^U \\
 \eta_\nu^V
 \end{vmatrix}$$

Jones & O'Dell (1977)

$$\begin{array}{c|c|c|c|c} \left(\frac{d}{dl} + \kappa_I \right) & \kappa_Q & \kappa_U & \kappa_V & I_\nu \\ \kappa_Q & \left(\frac{d}{dl} + \kappa_I \right) & \boxed{\kappa^*_V} & -\kappa^*_U & Q_\nu \\ \kappa_U & -\kappa^*_V & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_Q & U_\nu \\ \kappa_V & \kappa^*_U & -\kappa^*_Q & \left(\frac{d}{dl} + \kappa_I \right) & V_\nu \end{array} = \begin{array}{c} \eta_\nu^I \\ \eta_\nu^Q \\ \eta_\nu^U \\ \eta_\nu^V \end{array}$$

Jones & O'Dell (1977)

Faraday Rotation:



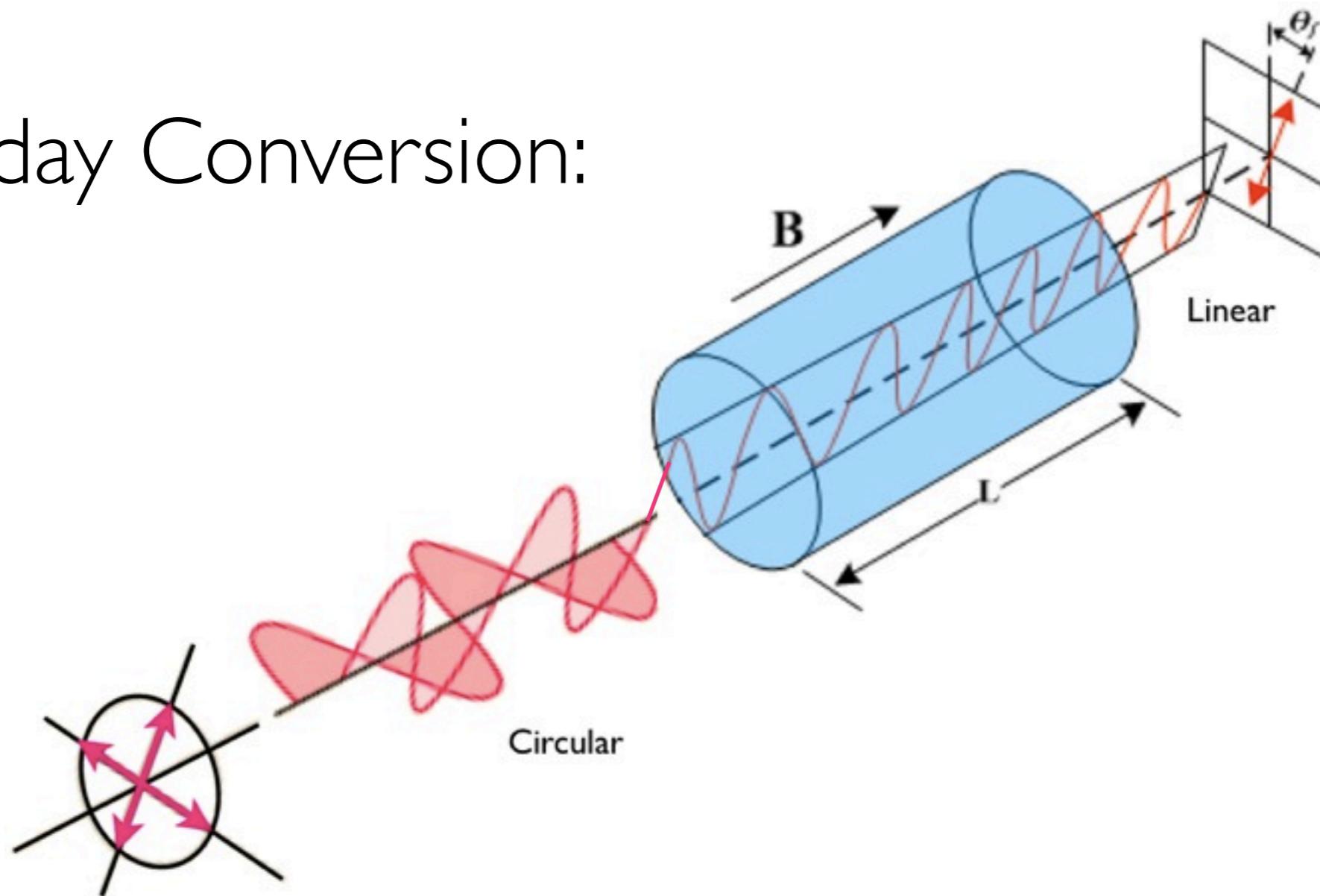
$$\begin{vmatrix}
\left(\frac{d}{dl} + \kappa_I \right) & \kappa_Q & \kappa_U & \kappa_V & I_\nu \\
\kappa_Q & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_V & -\kappa^*_U & Q_\nu \\
\kappa_U & -\kappa^*_V & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_Q & U_\nu \\
\kappa_V & \kappa^*_U & -\kappa^*_Q & \left(\frac{d}{dl} + \kappa_I \right) & V_\nu
\end{vmatrix} = \begin{vmatrix} \eta_\nu^I \\ \eta_\nu^Q \\ \eta_\nu^U \\ \eta_\nu^V \end{vmatrix}$$

Jones & O'Dell (1977)

$$\begin{array}{c|ccc|c}
& \left(\frac{d}{dl} + \kappa_I \right) & \kappa_Q & \kappa_U & \kappa_V \\
\kappa_Q & \left(\frac{d}{dl} + \kappa_I \right) & & & \\
& -\kappa^*_V & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_V & -\kappa^*_U \\
\kappa_U & & & & \kappa^*_Q \\
& \boxed{\kappa^*_U} & & \boxed{-\kappa^*_Q} & \left(\frac{d}{dl} + \kappa_I \right) \\
\kappa_V & & & &
\end{array} \quad \begin{array}{c|c|c}
I_\nu & Q_\nu & \eta_\nu^I \\
U_\nu & U_\nu & \eta_\nu^Q \\
V_\nu & V_\nu & \eta_\nu^V
\end{array} = \begin{array}{c}
\end{array}$$

Jones & O'Dell (1977)

Faraday Conversion:



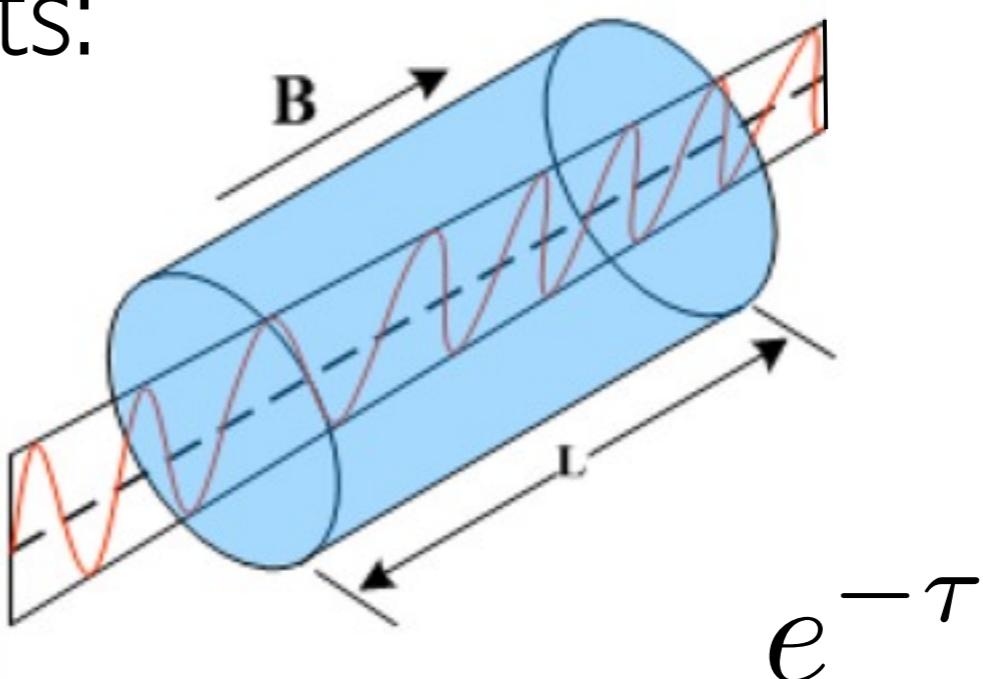
$$\begin{vmatrix}
\left(\frac{d}{dl} + \kappa_I \right) & \kappa_Q & \kappa_U & \kappa_V & I_\nu \\
\kappa_Q & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_V & -\kappa^*_U & Q_\nu \\
\kappa_U & -\kappa^*_V & \left(\frac{d}{dl} + \kappa_I \right) & \kappa^*_Q & U_\nu \\
\kappa_V & \kappa^*_U & -\kappa^*_Q & \left(\frac{d}{dl} + \kappa_I \right) & V_\nu
\end{vmatrix} = \begin{vmatrix} \eta_\nu^I \\ \eta_\nu^Q \\ \eta_\nu^U \\ \eta_\nu^V \end{vmatrix}$$

Jones & O'Dell (1977)

$$\begin{array}{c}
\left(\frac{d}{dl} + \kappa_I \right) \\
\kappa_Q \\
\kappa_U \\
\kappa_V
\end{array}
\quad
\begin{array}{c}
\kappa_Q \\
\left(\frac{d}{dl} + \kappa_I \right) \\
-\kappa^*_V \\
\kappa^*_U
\end{array}
\quad
\begin{array}{c}
\kappa_U \\
\kappa^*_V \\
\left(\frac{d}{dl} + \kappa_I \right) \\
-\kappa^*_Q
\end{array}
\quad
\begin{array}{c}
\kappa_V \\
-\kappa^*_U \\
\kappa^*_Q \\
\left(\frac{d}{dl} + \kappa_I \right)
\end{array}
\quad
\begin{array}{c}
I_\nu \\
Q_\nu \\
U_\nu \\
V_\nu
\end{array}
=
\begin{array}{c}
\eta_\nu^I \\
\eta_\nu^Q \\
\eta_\nu^U \\
\eta_\nu^V
\end{array}$$

Jones & O'Dell (1977)

Optical Depth Effects:



The Full Stokes Equations for Radiative Transfer of Polarized Emission:

$$\begin{vmatrix} I_\nu \\ Q_\nu \\ U_\nu \\ V_\nu \end{vmatrix} = \begin{vmatrix} I_\nu^\infty \\ Q_\nu^\infty \\ U_\nu^\infty \\ V_\nu^\infty \end{vmatrix} + e^{-\tau} \left\{ \begin{array}{c} \text{A lot of math!} \\ (B, \theta, L, \nu, n_o, \alpha, \gamma_{\min}) \end{array} \right\} \begin{vmatrix} I_\nu^0 - I_\nu^\infty \\ Q_\nu^0 - Q_\nu^\infty \\ U_\nu^0 - U_\nu^\infty \\ V_\nu^0 - V_\nu^\infty \end{vmatrix}$$

Jones & O'Dell (1977)

The Full Stokes Equations for Radiative Transfer of Polarized Emission:

$$\begin{vmatrix} I_\nu \\ Q_\nu \\ U_\nu \\ V_\nu \end{vmatrix} = \begin{vmatrix} I_\nu^\infty \\ Q_\nu^\infty \\ U_\nu^\infty \\ V_\nu^\infty \end{vmatrix} + e^{-\tau} \left\{ \begin{array}{c} \text{A lot of math!} \\ (B, \theta, L, \nu, n_o, \alpha, \gamma_{\min}) \end{array} \right\} \begin{vmatrix} I_\nu^0 - I_\nu^\infty \\ Q_\nu^0 - Q_\nu^\infty \\ U_\nu^0 - U_\nu^\infty \\ V_\nu^0 - V_\nu^\infty \end{vmatrix}$$

Jones & O'Dell (1977)

The Full Stokes Equations for Radiative Transfer of Polarized Emission:

$$\begin{vmatrix} I_\nu \\ Q_\nu \\ U_\nu \\ V_\nu \end{vmatrix} = \begin{vmatrix} I_\nu^\infty \\ Q_\nu^\infty \\ U_\nu^\infty \\ V_\nu^\infty \end{vmatrix} + e^{-\tau} \left\{ \begin{array}{c} \\ \\ \\ \end{array} \right. \quad \left. \begin{array}{c} \text{A lot of math!} \\ (B, \theta, L, \nu, n_o, \alpha, \gamma_{\min}) \end{array} \right\} \begin{vmatrix} I_\nu^0 - I_\nu^\infty \\ Q_\nu^0 - Q_\nu^\infty \\ U_\nu^0 - U_\nu^\infty \\ V_\nu^0 - V_\nu^\infty \end{vmatrix}$$

Jones & O'Dell (1977)

The Full Stokes Equations for Radiative Transfer of Polarized Emission:

$$\begin{array}{c|c} I_\nu & I_\nu^\infty \\ \hline Q_\nu & Q_\nu^\infty \\ \hline U_\nu & U_\nu^\infty \\ \hline V_\nu & V_\nu^\infty \end{array} = \begin{array}{c|c} I_\nu^\infty \\ \hline Q_\nu^\infty \\ \hline U_\nu^\infty \\ \hline V_\nu^\infty \end{array} + e^{-\tau} \left\{ \begin{array}{c} \text{A lot of math!} \\ (B, \theta, L, \nu, n_o, \alpha, \gamma_{\min}) \end{array} \right\} \begin{array}{c|c} I_\nu^0 - I_\nu^\infty \\ \hline Q_\nu^0 - Q_\nu^\infty \\ \hline U_\nu^0 - U_\nu^\infty \\ \hline V_\nu^0 - V_\nu^\infty \end{array}$$

Jones & O'Dell (1977)

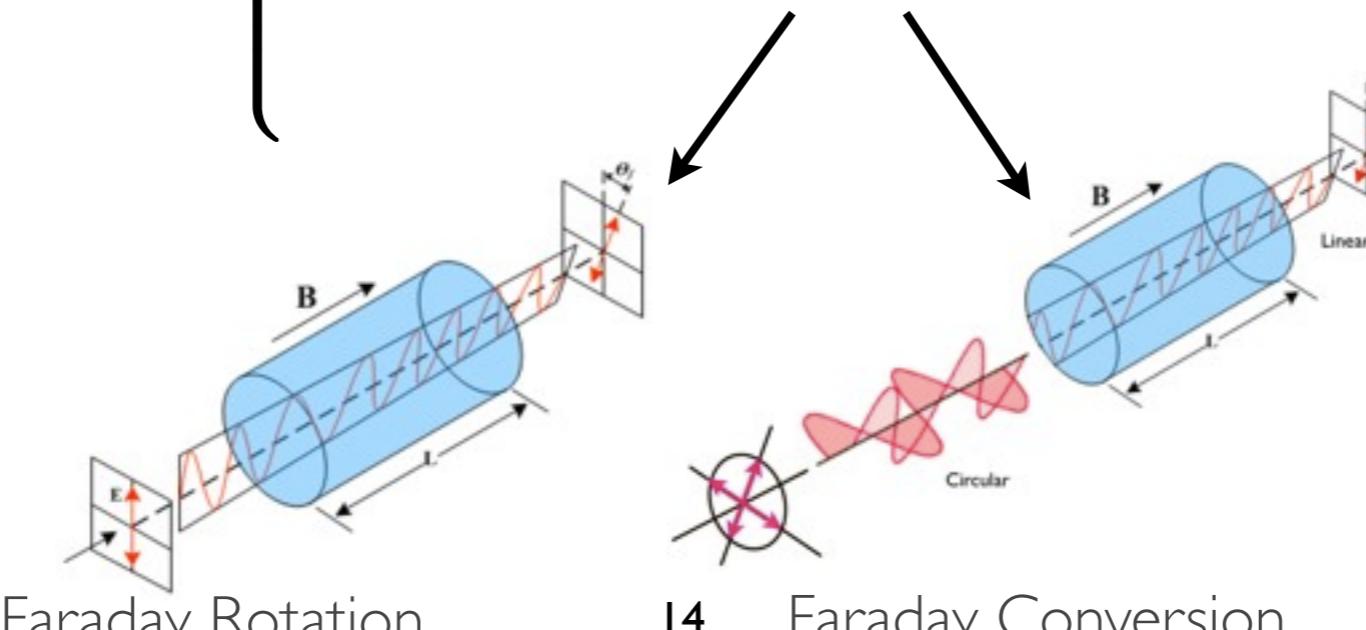
The Full Stokes Equations for Radiative Transfer of Polarized Emission:

$$\begin{vmatrix} I_\nu \\ Q_\nu \\ U_\nu \\ V_\nu \end{vmatrix} = \begin{vmatrix} I_\nu^\infty \\ Q_\nu^\infty \\ U_\nu^\infty \\ V_\nu^\infty \end{vmatrix} + e^{-\tau} \left\{ \begin{array}{c} \text{A lot of math!} \\ (B, \theta, L, \nu, n_o, \alpha, \gamma_{\min}) \end{array} \right\} \begin{vmatrix} I_\nu^0 - I_\nu^\infty \\ Q_\nu^0 - Q_\nu^\infty \\ U_\nu^0 - U_\nu^\infty \\ V_\nu^0 - V_\nu^\infty \end{vmatrix}$$

Jones & O'Dell (1977)

The Full Stokes Equations for Radiative Transfer of Polarized Emission:

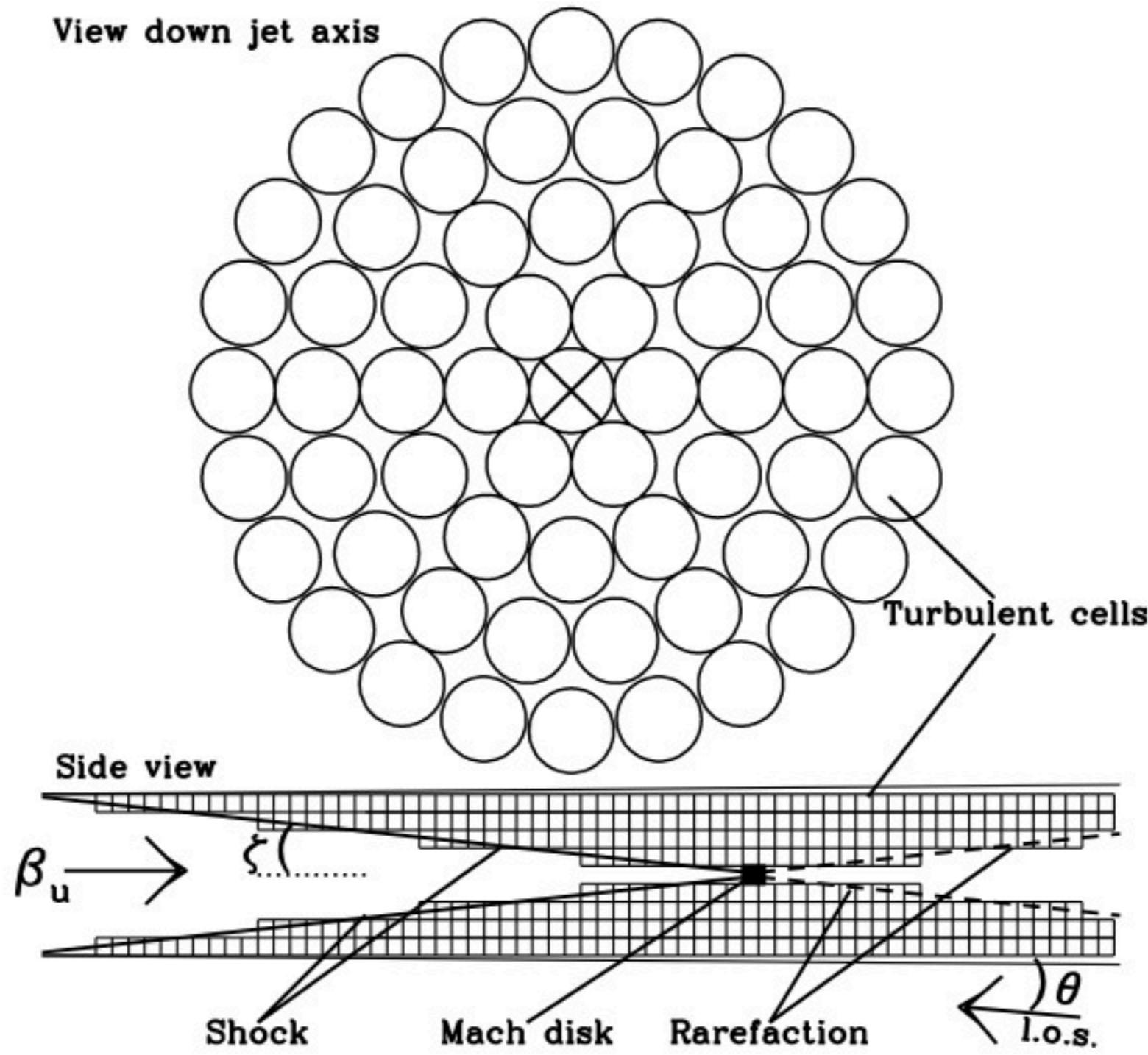
$$\begin{vmatrix} I_\nu \\ Q_\nu \\ U_\nu \\ V_\nu \end{vmatrix} = \begin{vmatrix} I_\nu^\infty \\ Q_\nu^\infty \\ U_\nu^\infty \\ V_\nu^\infty \end{vmatrix} + e^{-\tau} \left\{ \begin{array}{c} \text{A lot of math!} \\ (B, \theta, L, \nu, n_o, \alpha, \gamma_{\min}) \end{array} \right\} \begin{vmatrix} I_\nu^0 - I_\nu^\infty \\ Q_\nu^0 - Q_\nu^\infty \\ U_\nu^0 - U_\nu^\infty \\ V_\nu^0 - V_\nu^\infty \end{vmatrix}$$



 Faraday Rotation 14 Faraday Conversion

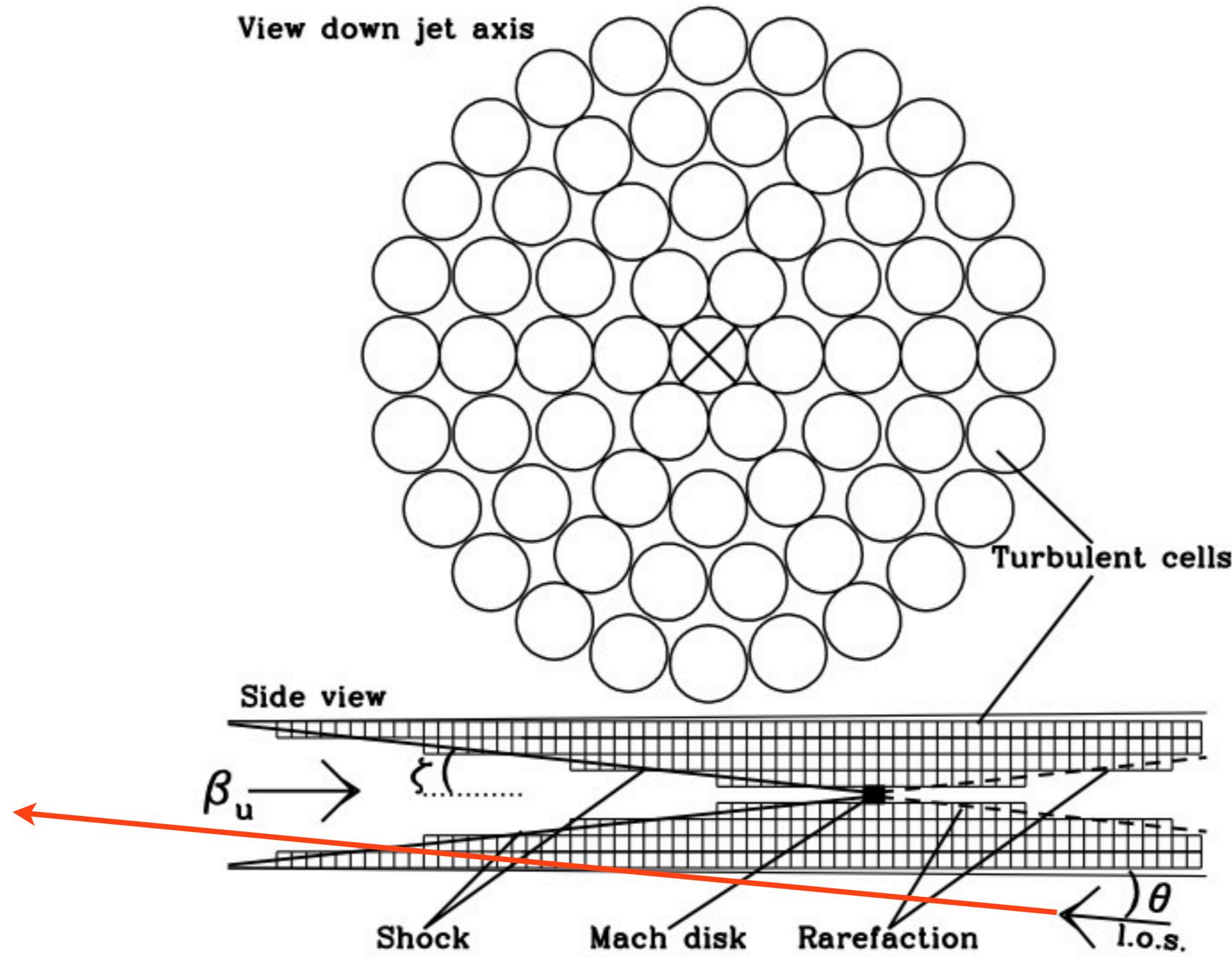
Jones & O'Dell (1977)

The turbulent extreme multi-zone (TEMZ) code:



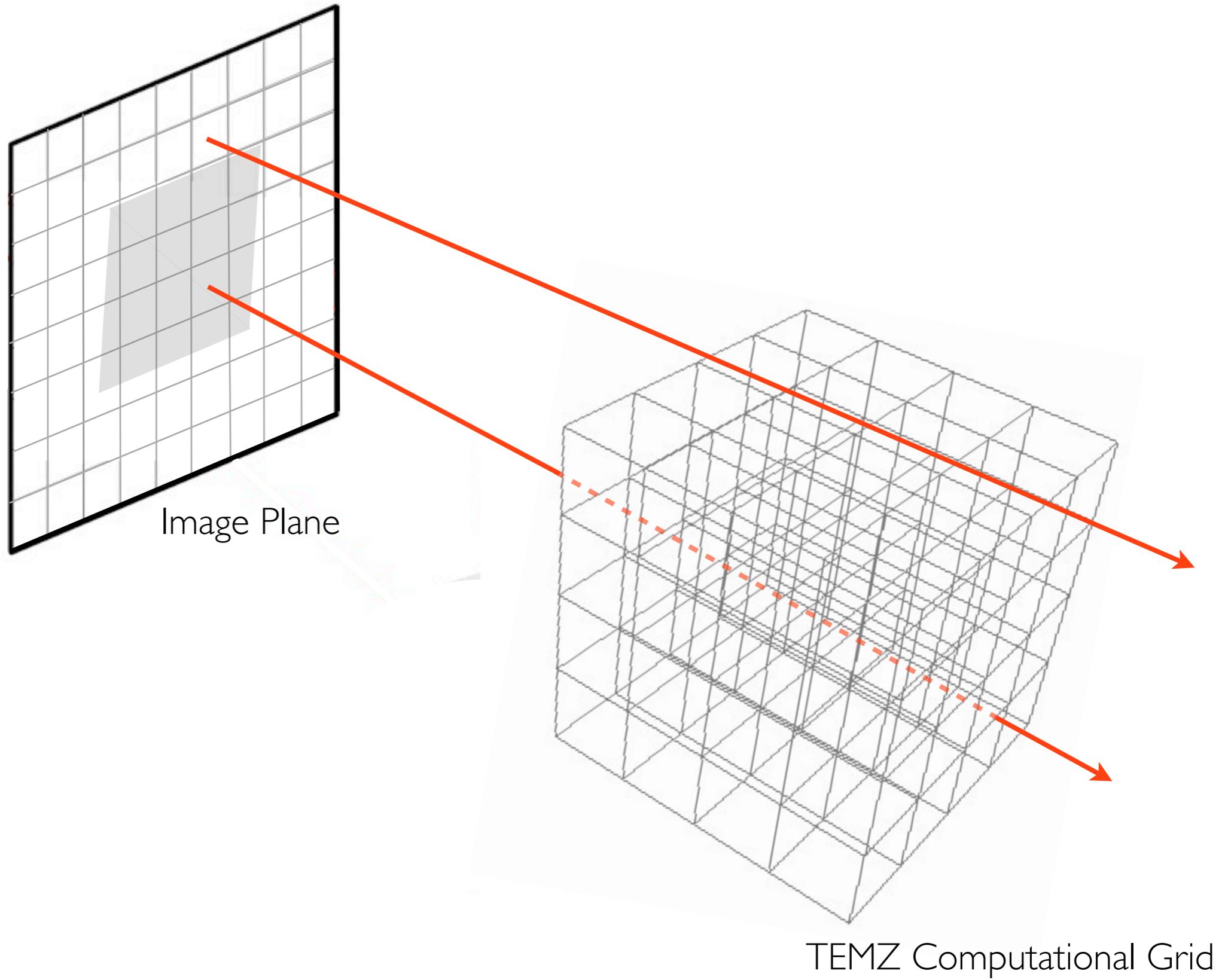
Marscher (2014)

The turbulent extreme multi-zone (TEMZ) code:

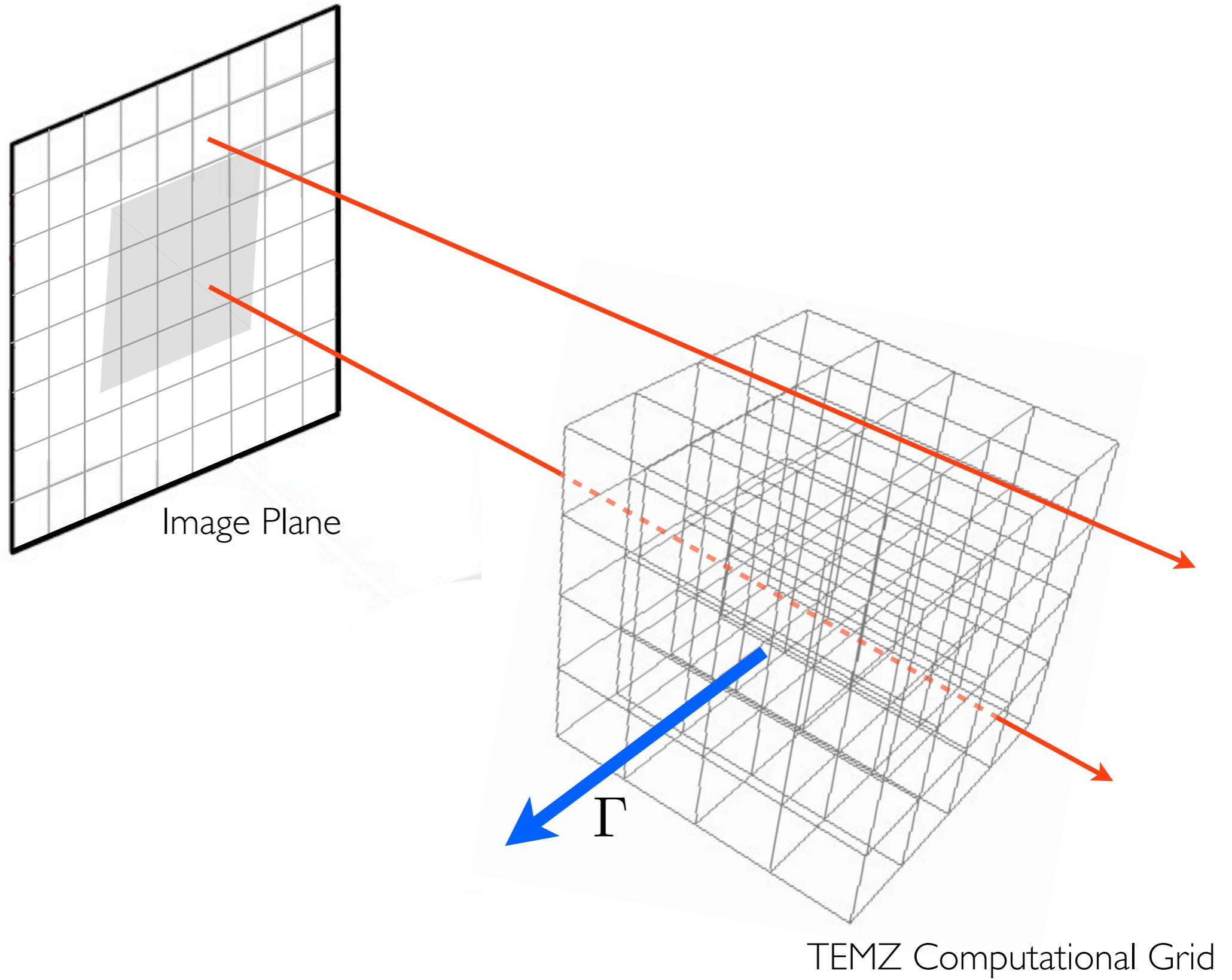


Marscher (2014)

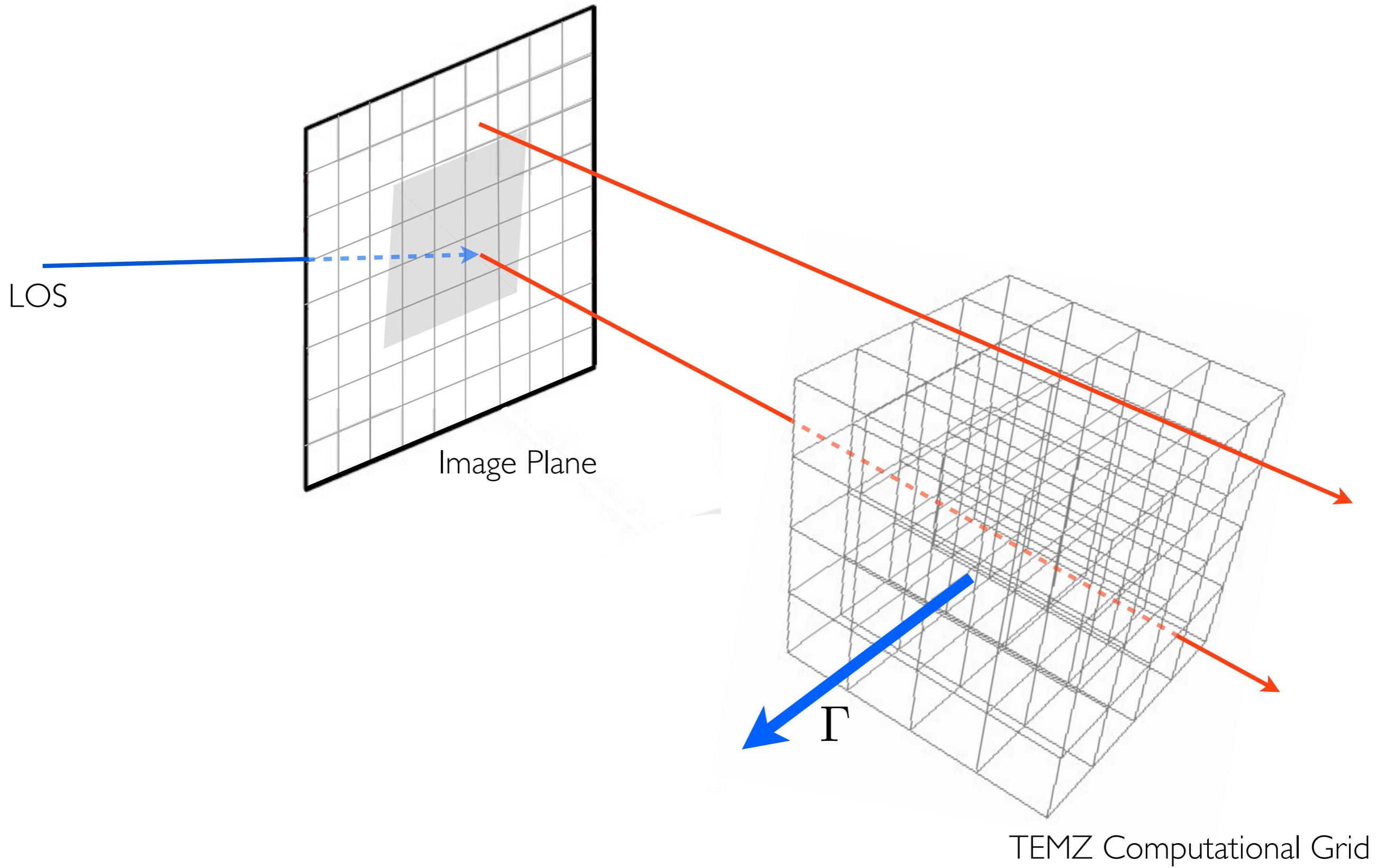
Ray Tracing Algorithm:



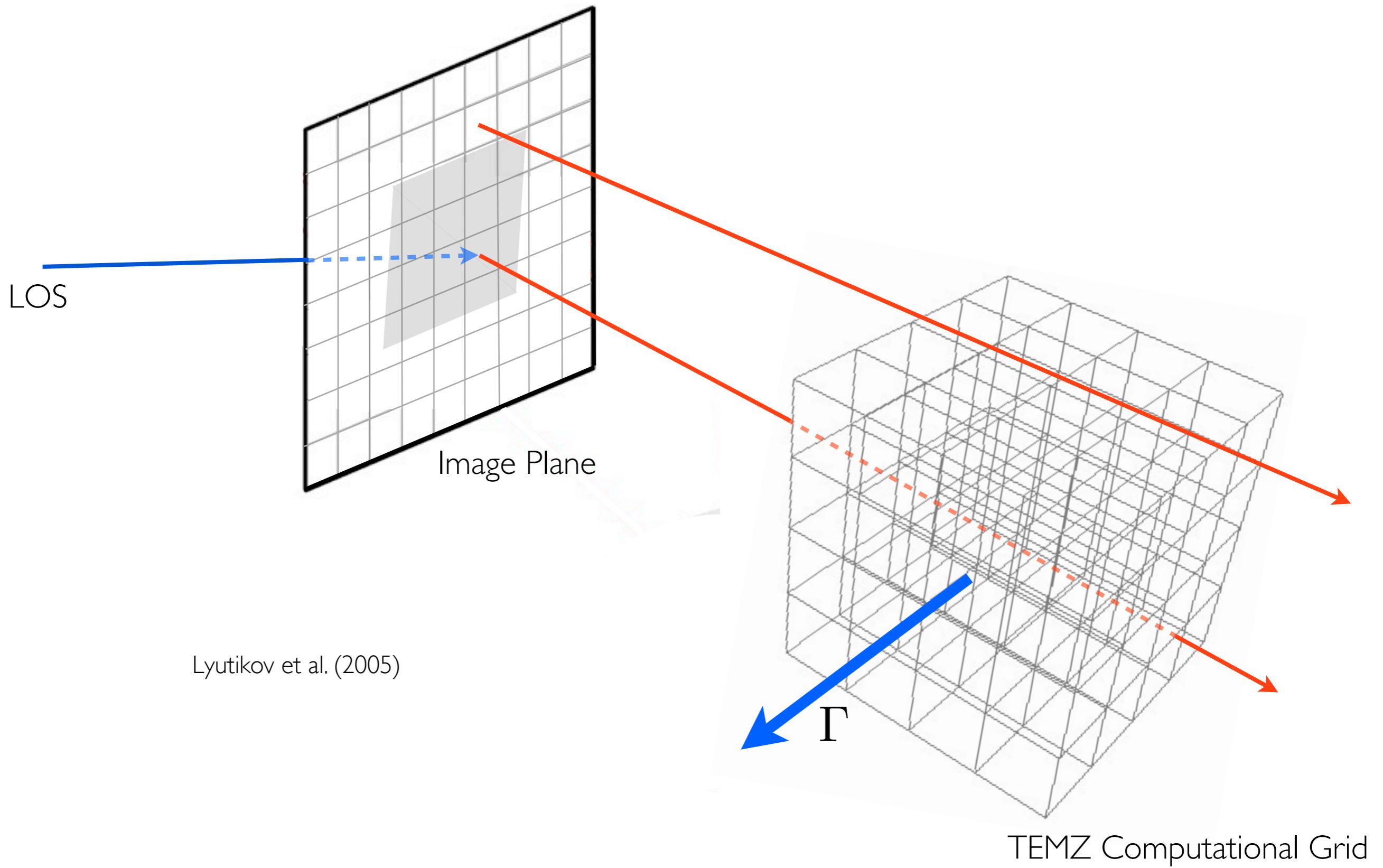
Ray Tracing Algorithm:



Ray Tracing Algorithm:



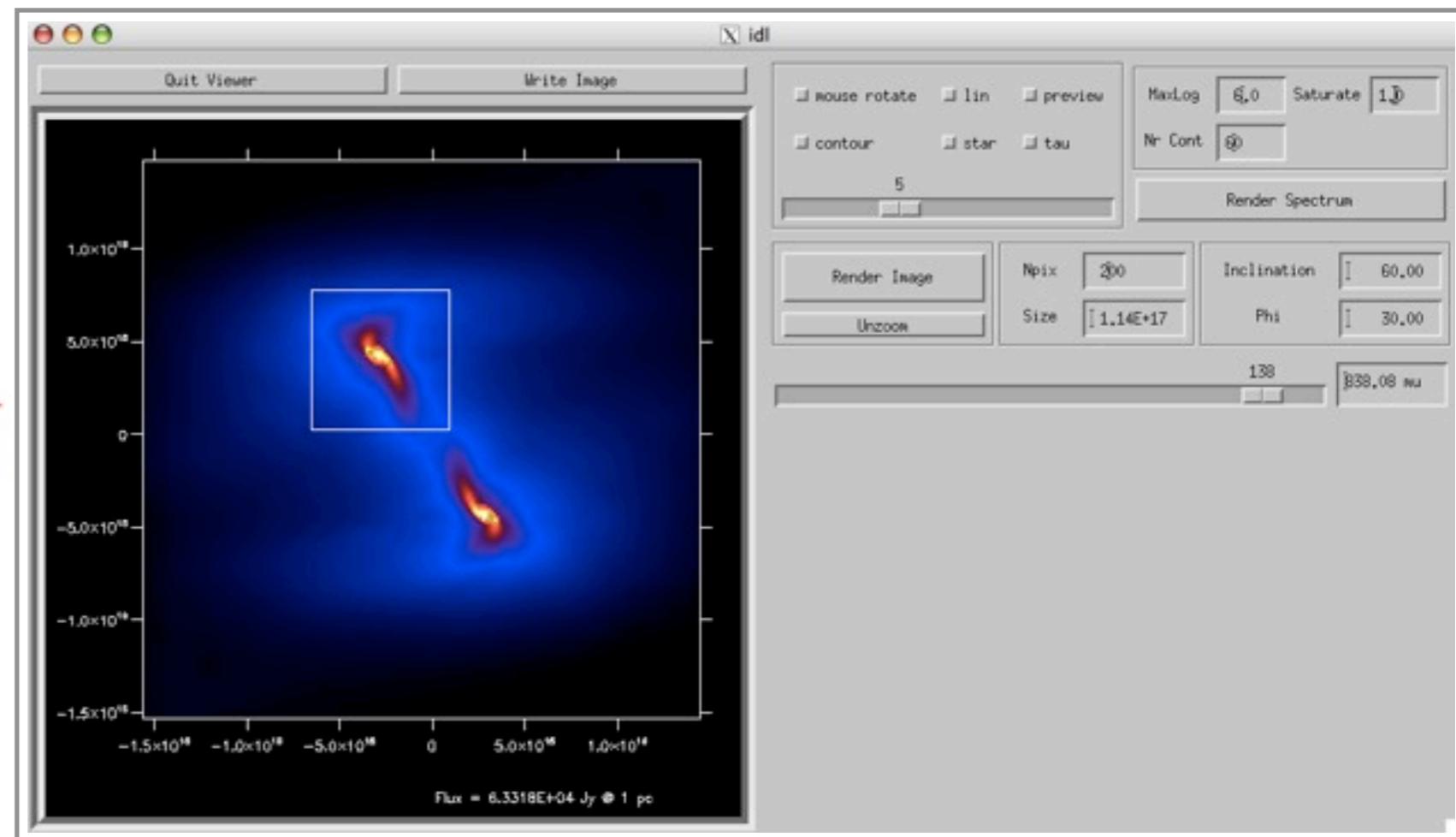
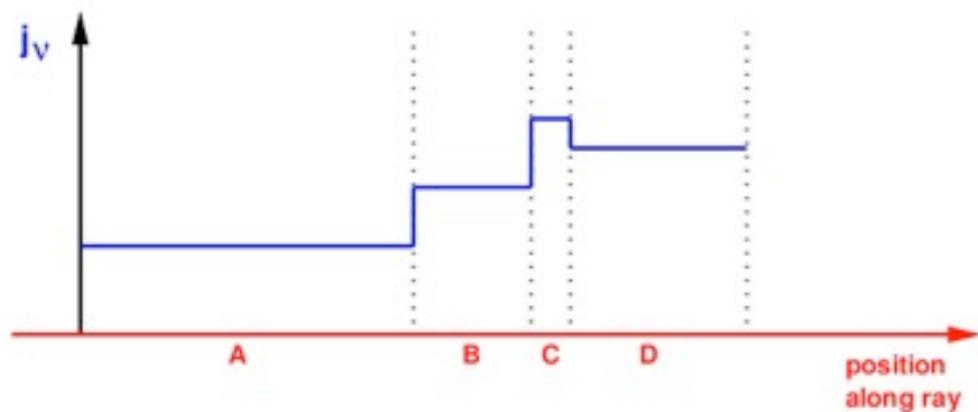
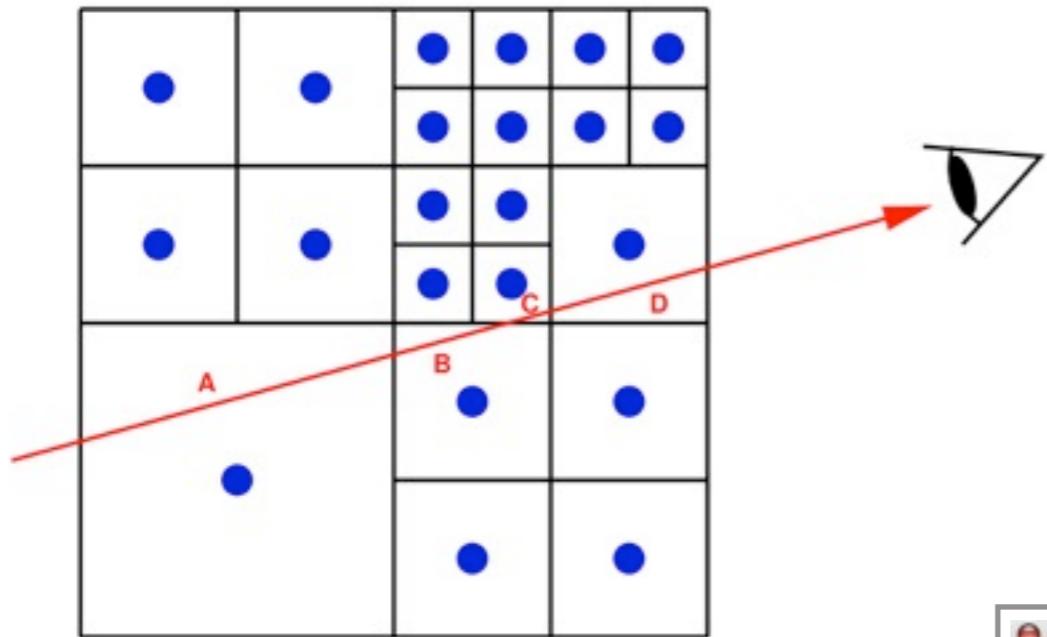
Ray Tracing Algorithm:



The RADMC-3D Code:



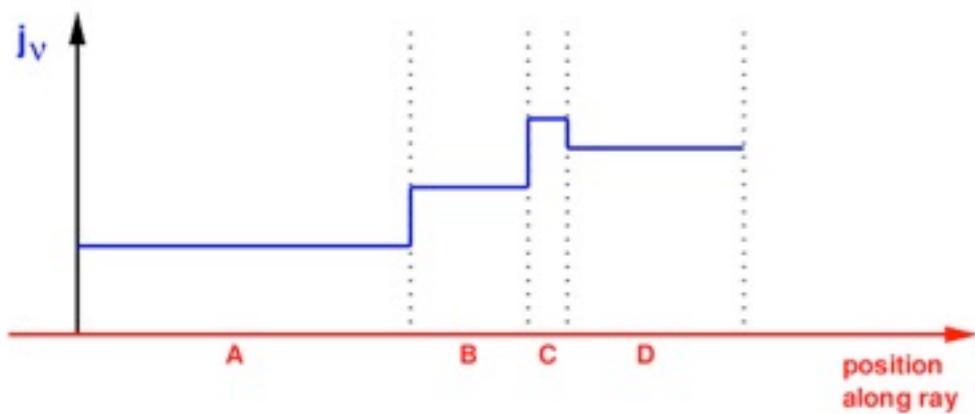
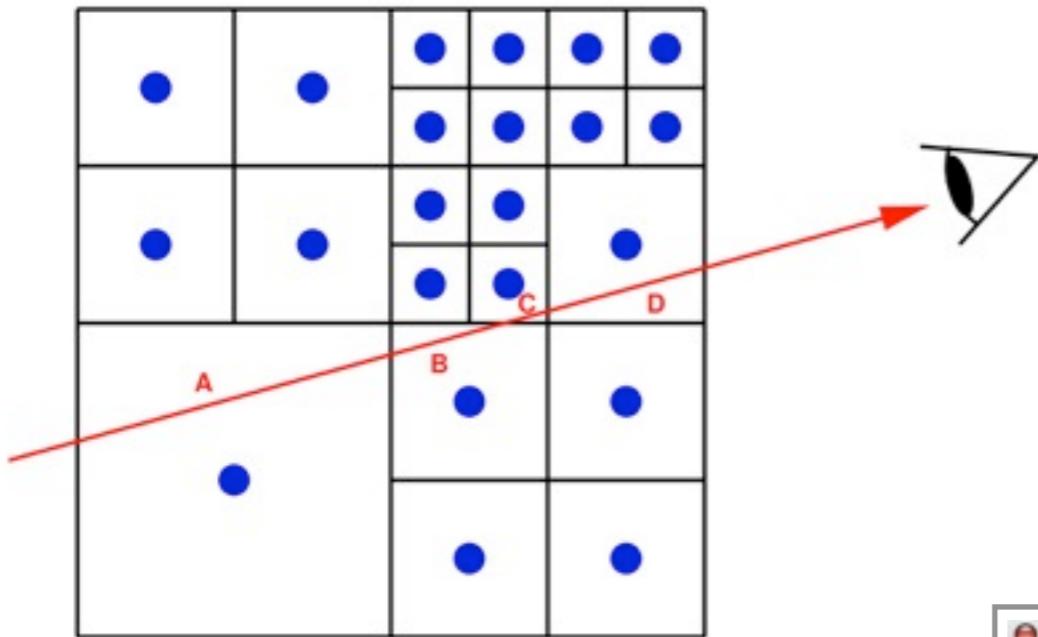
MAX-PLANCK-GESELLSCHAFT



The RADMC-3D Code:

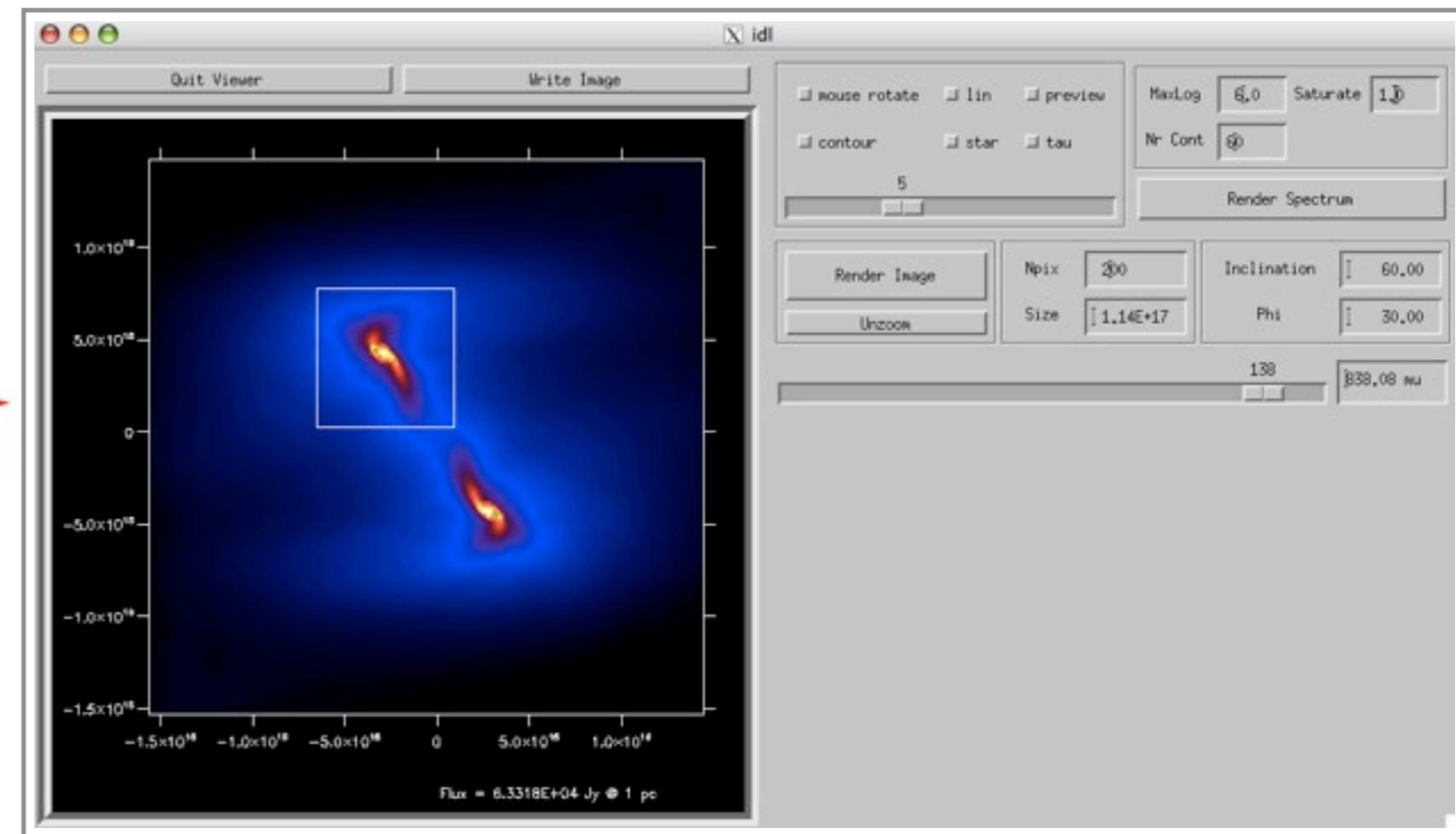


MAX-PLANCK-GESELLSCHAFT

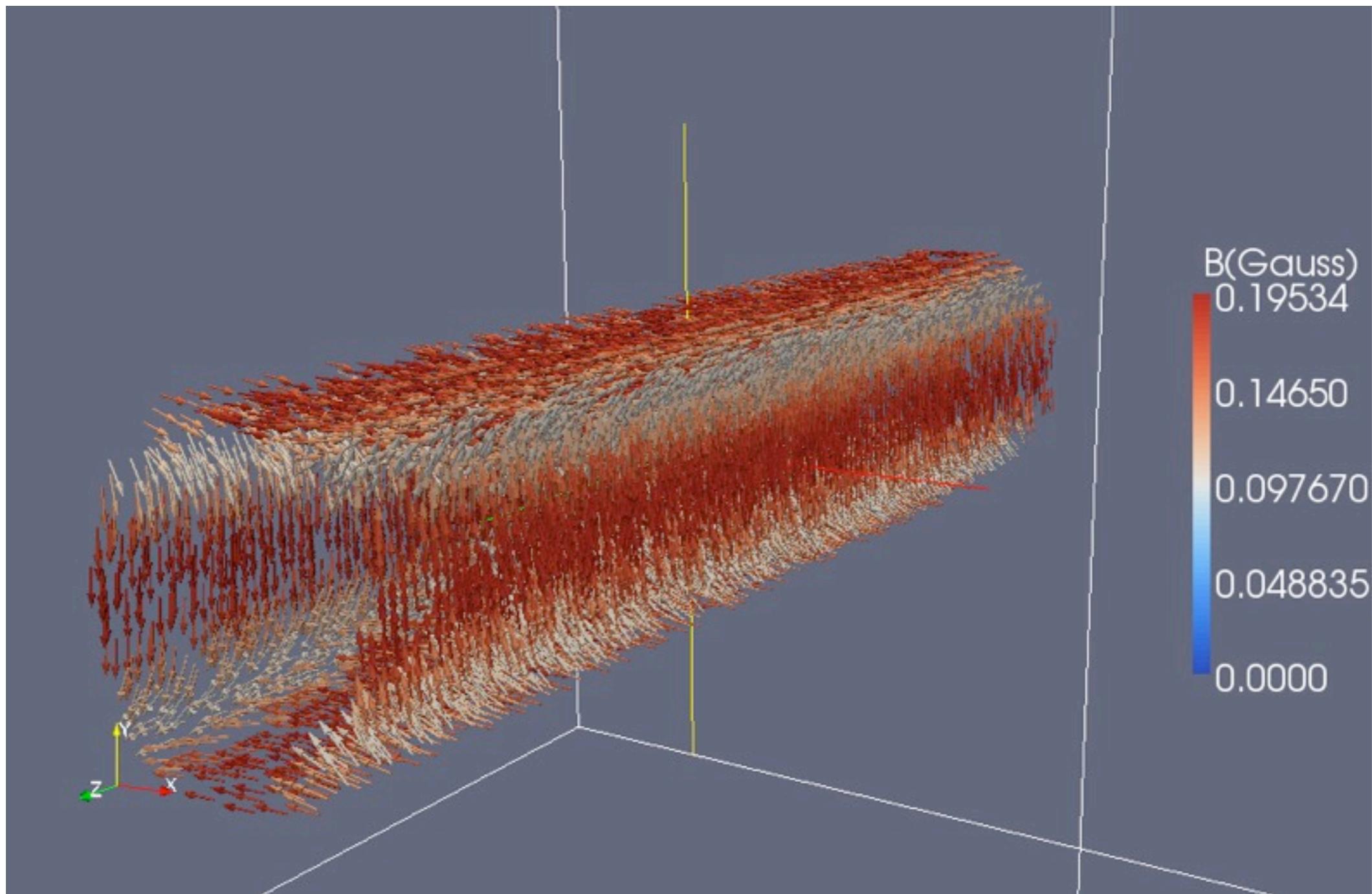


radmc3dPy

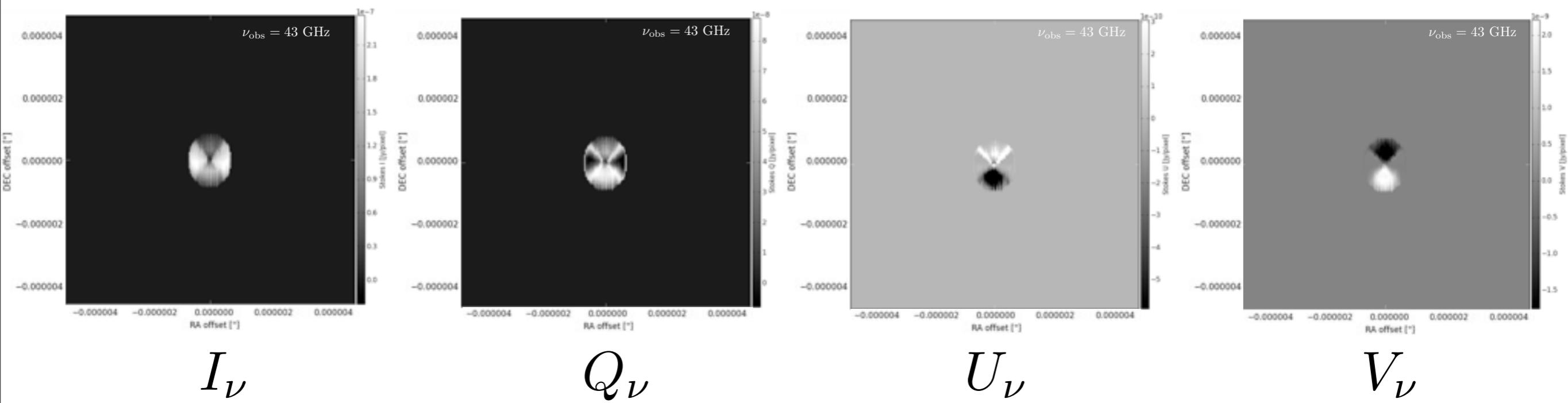
Attila Juhasz



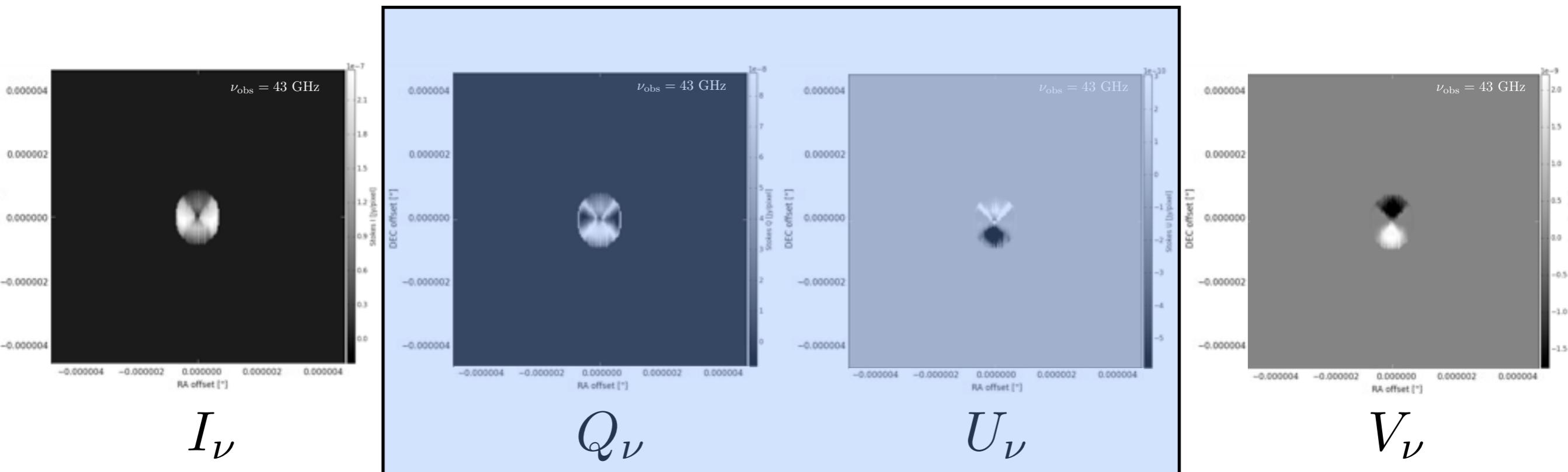
TEMZ Model (Order):



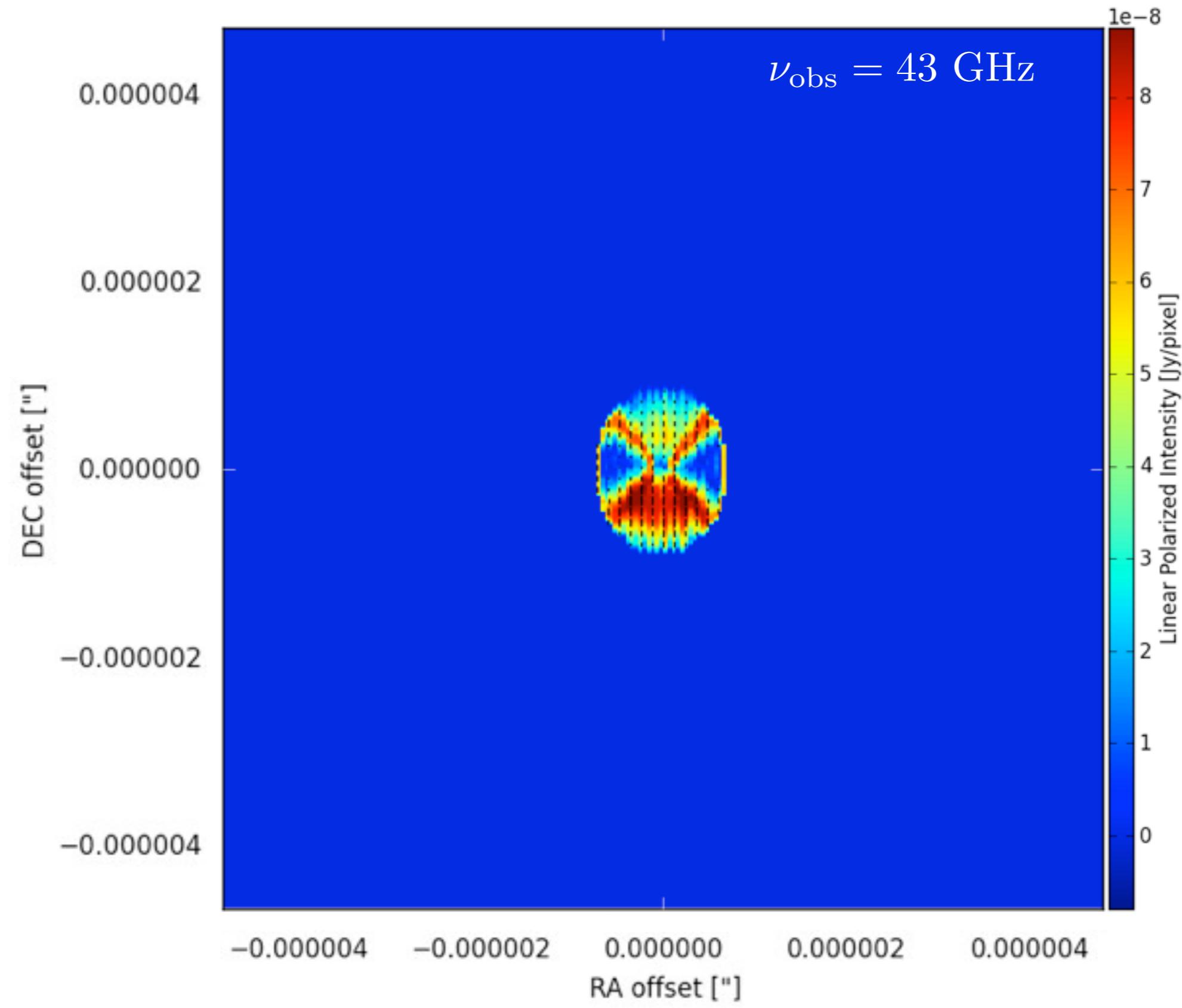
TEMZ Model (Order):



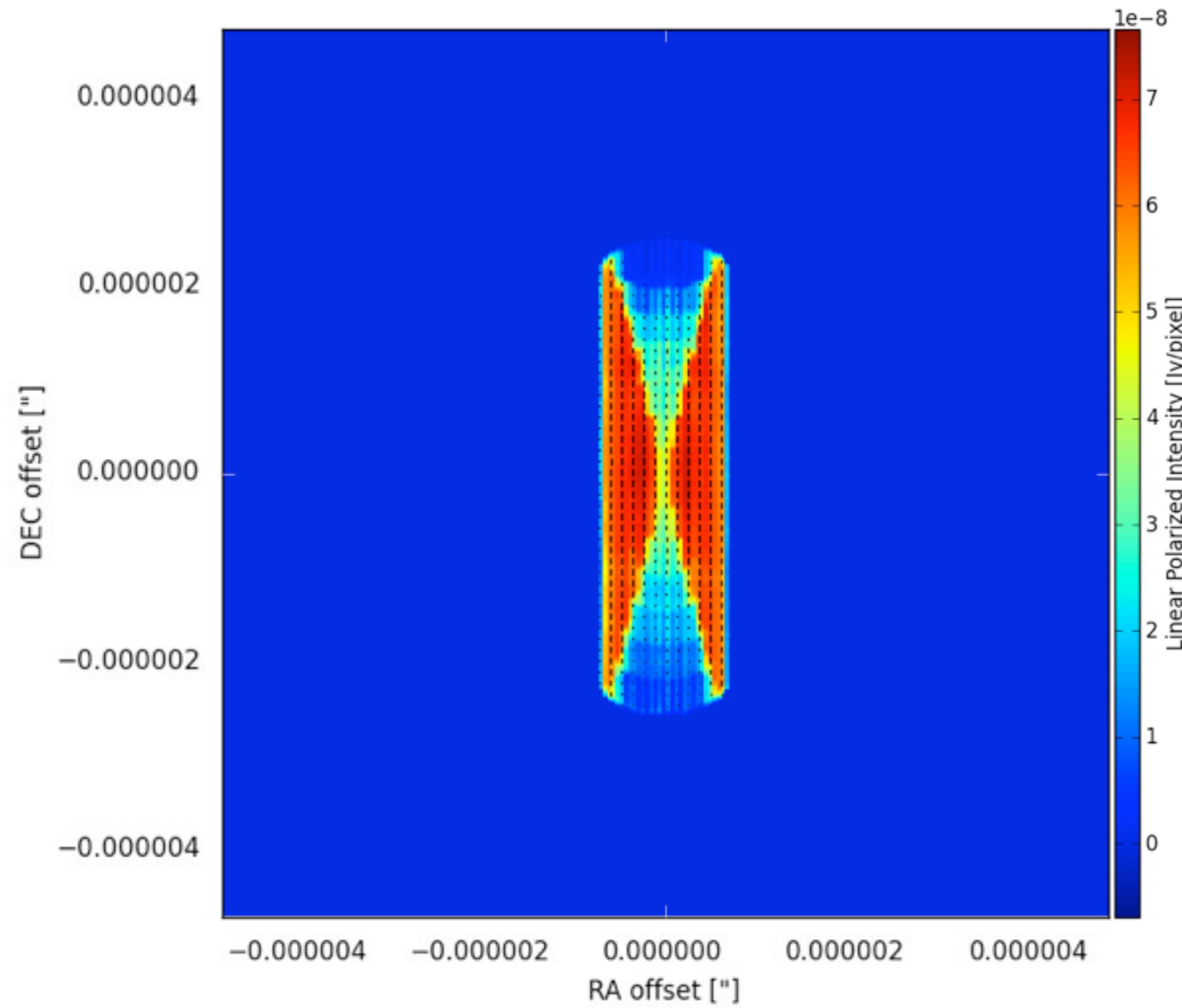
TEMZ Model (Order):



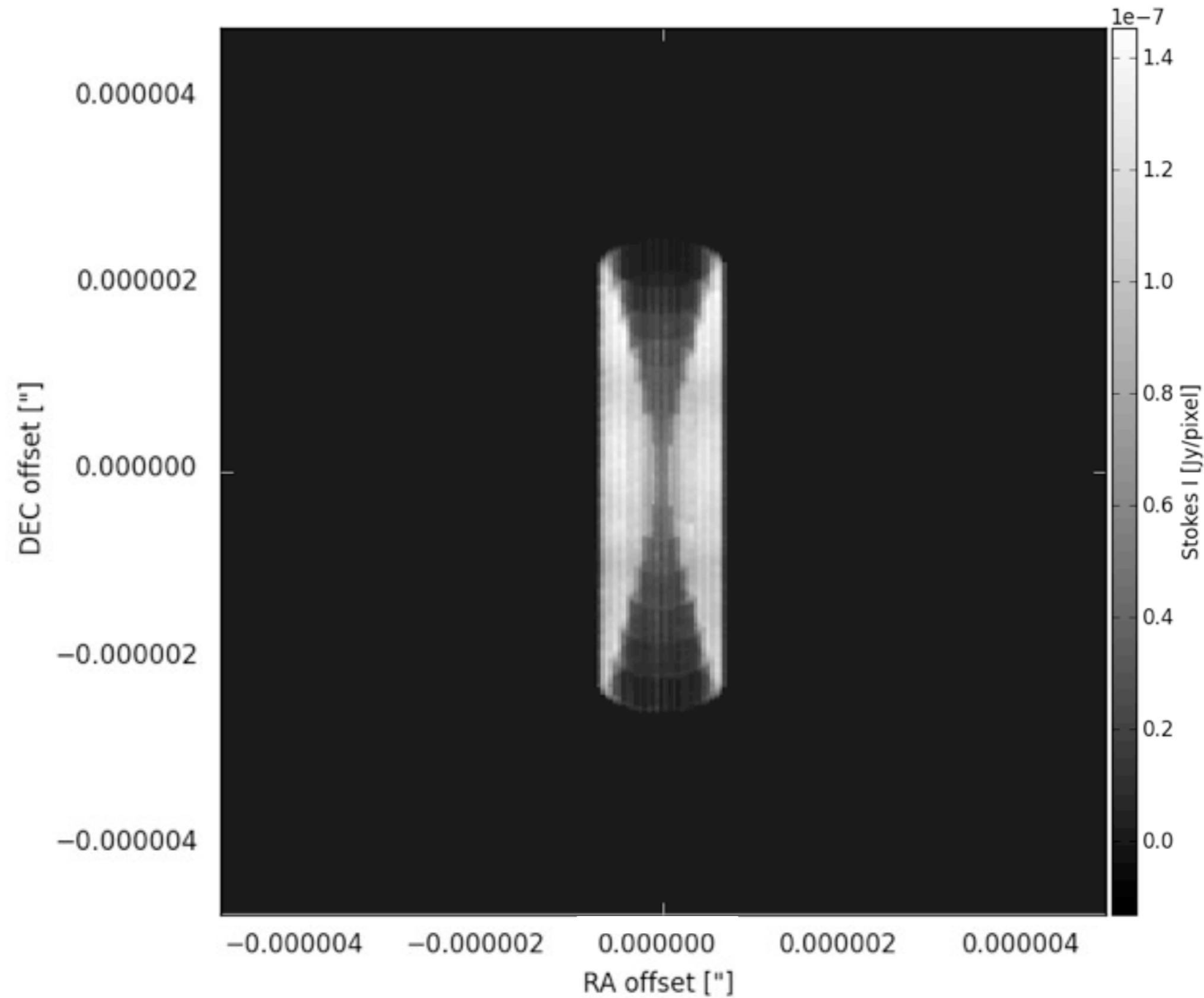
TEMZ Model (Order):



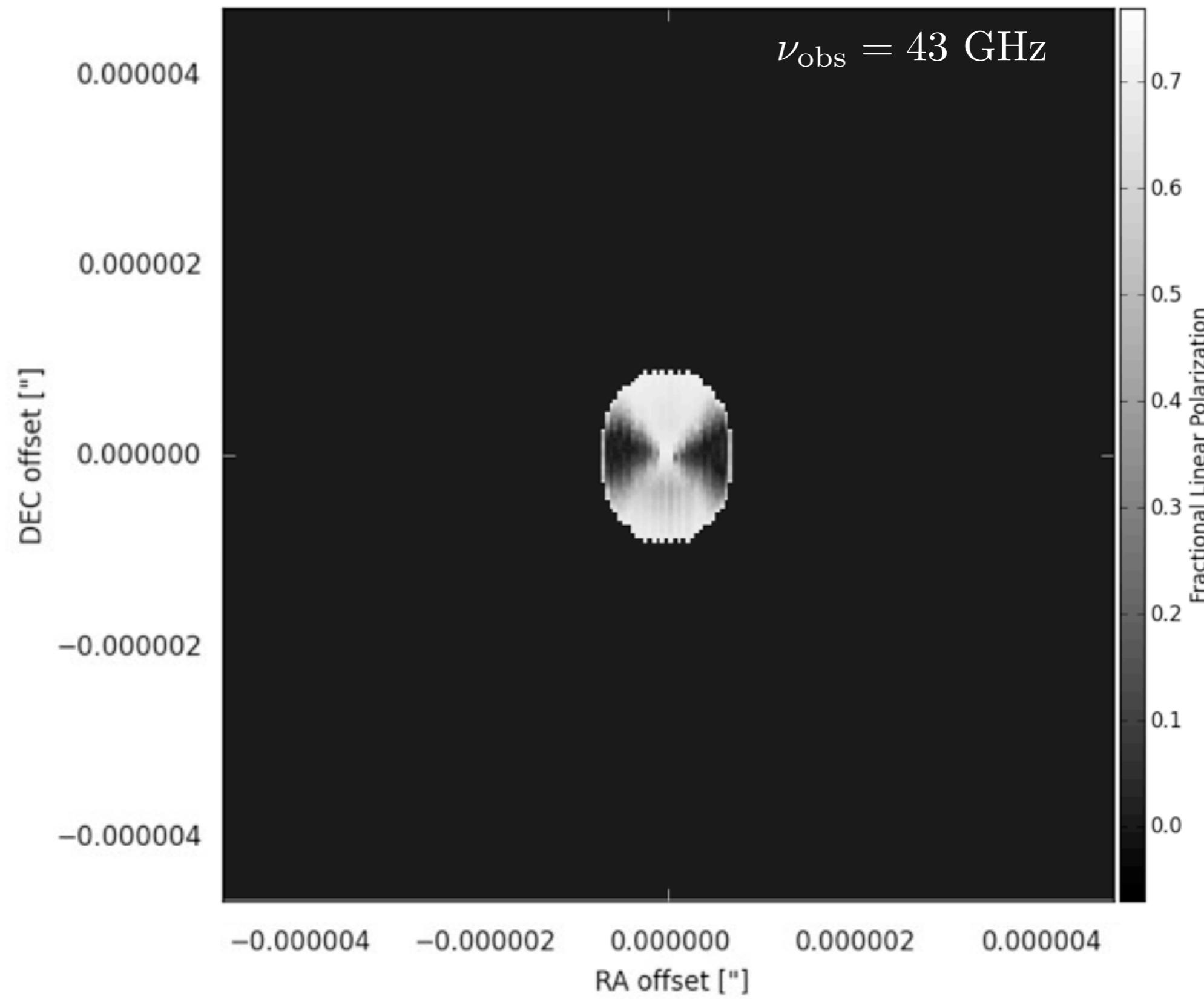
TEMZ Model (Order):



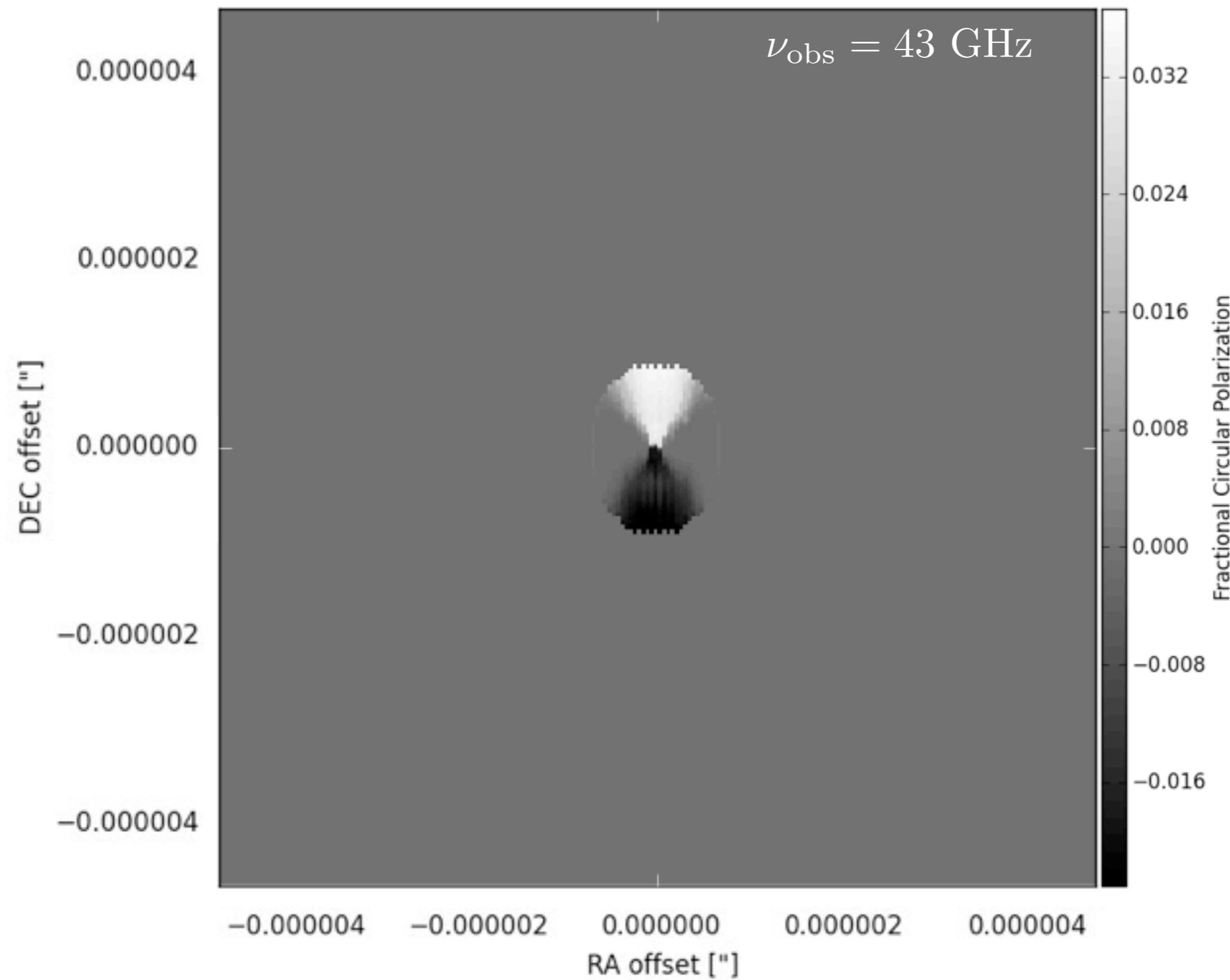
TEMZ Model (Order):



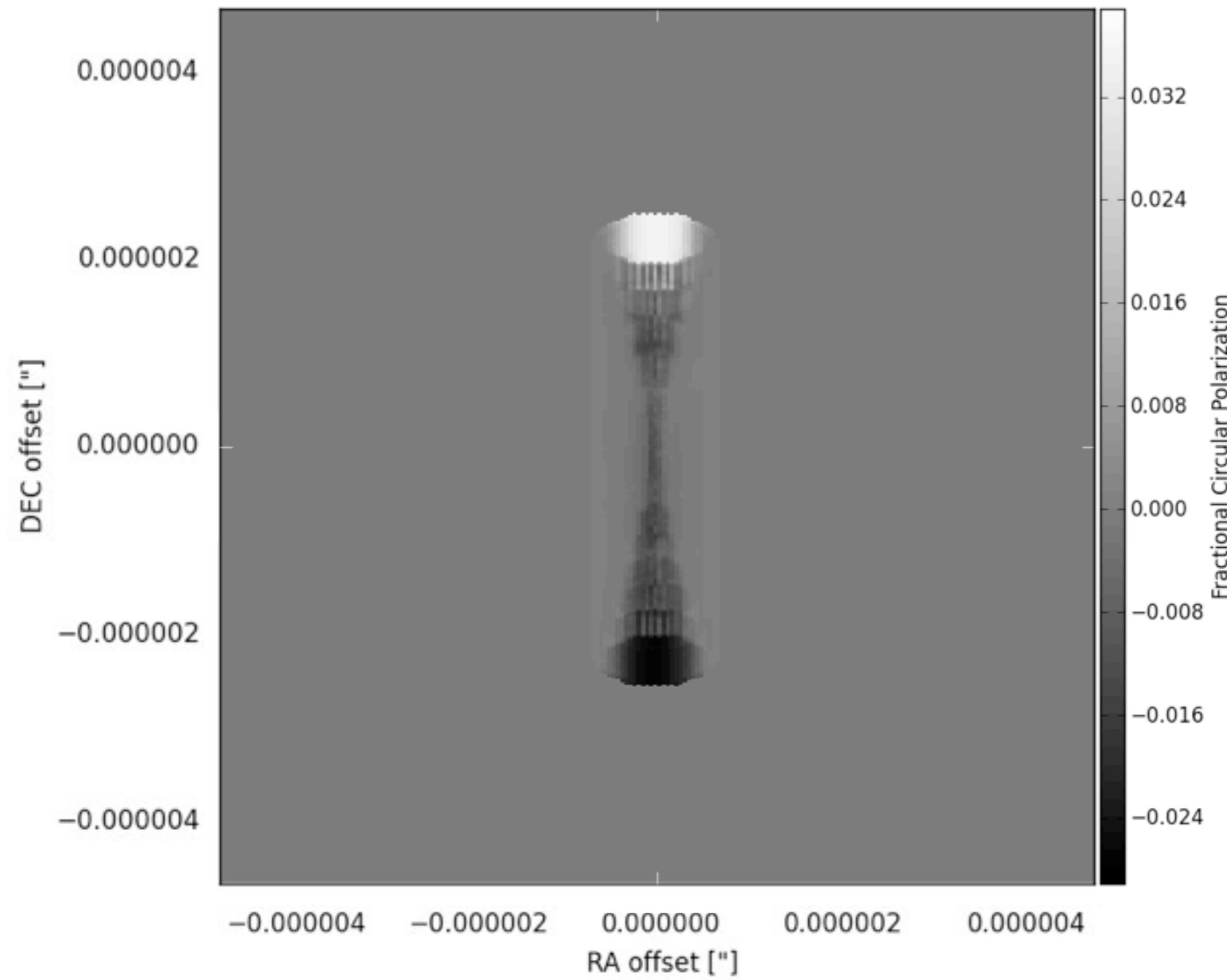
TEMZ Model (Order):



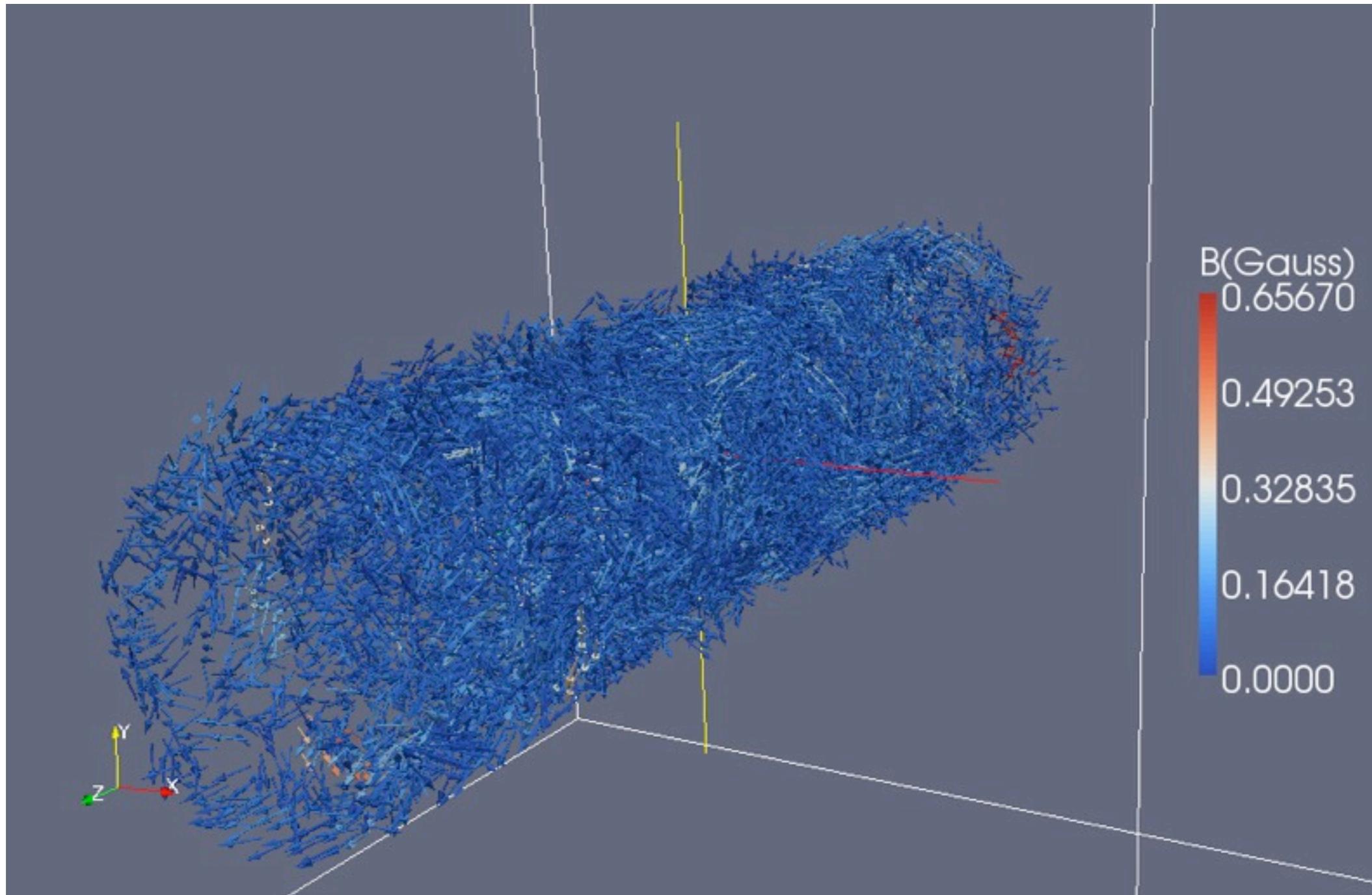
TEMZ Model (Order):



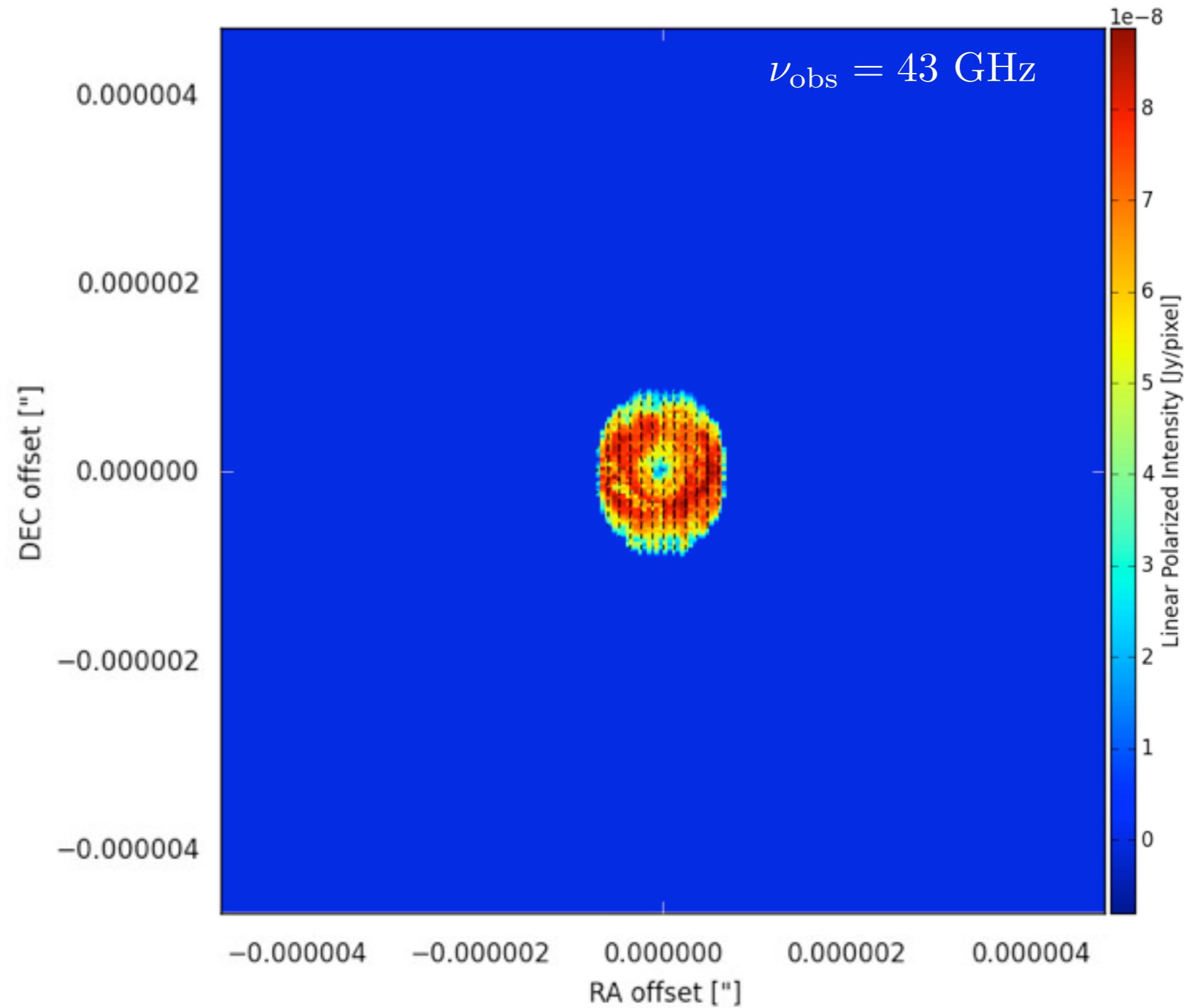
TEMZ Model (Order):



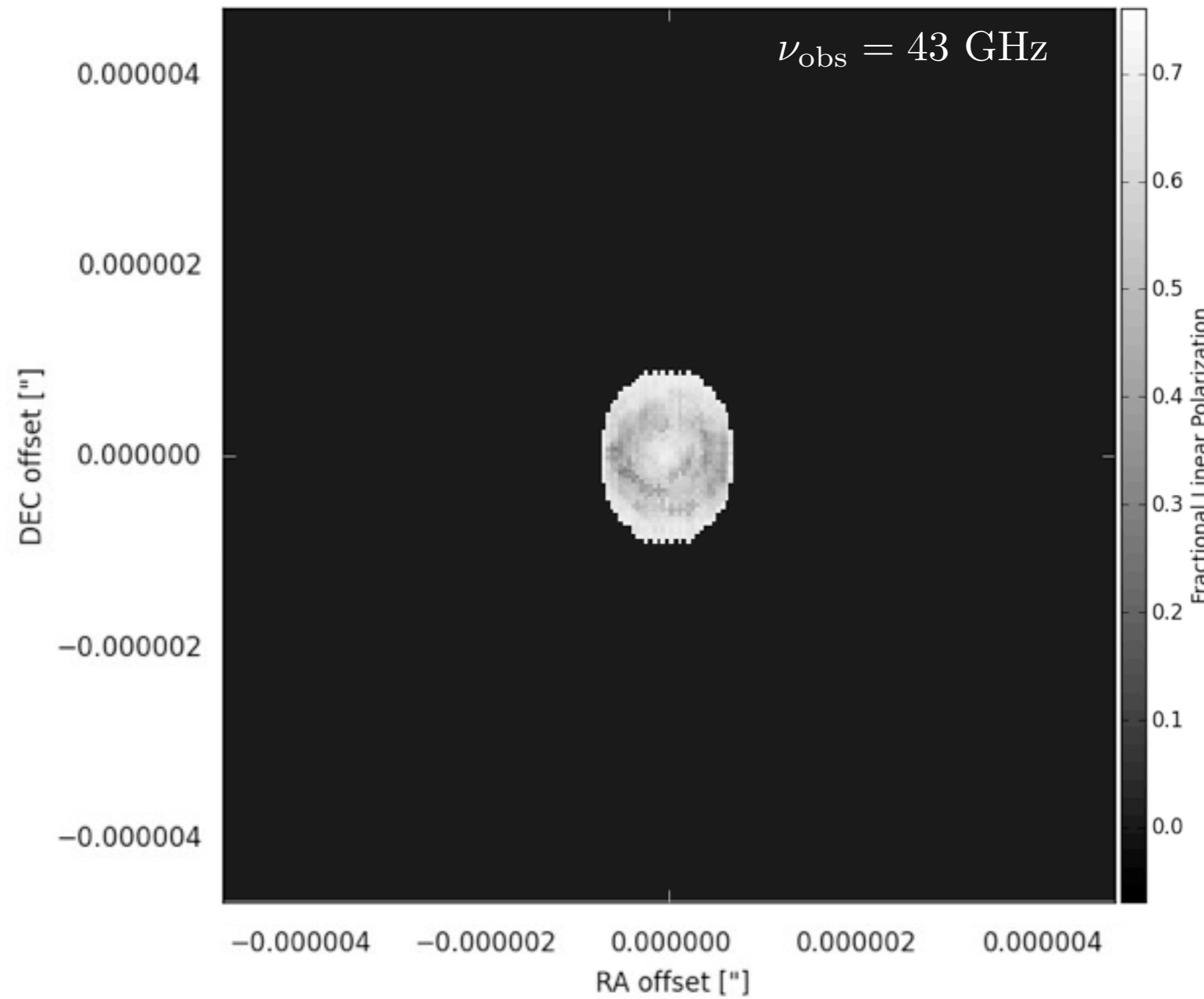
TEMZ Model (Disorder):



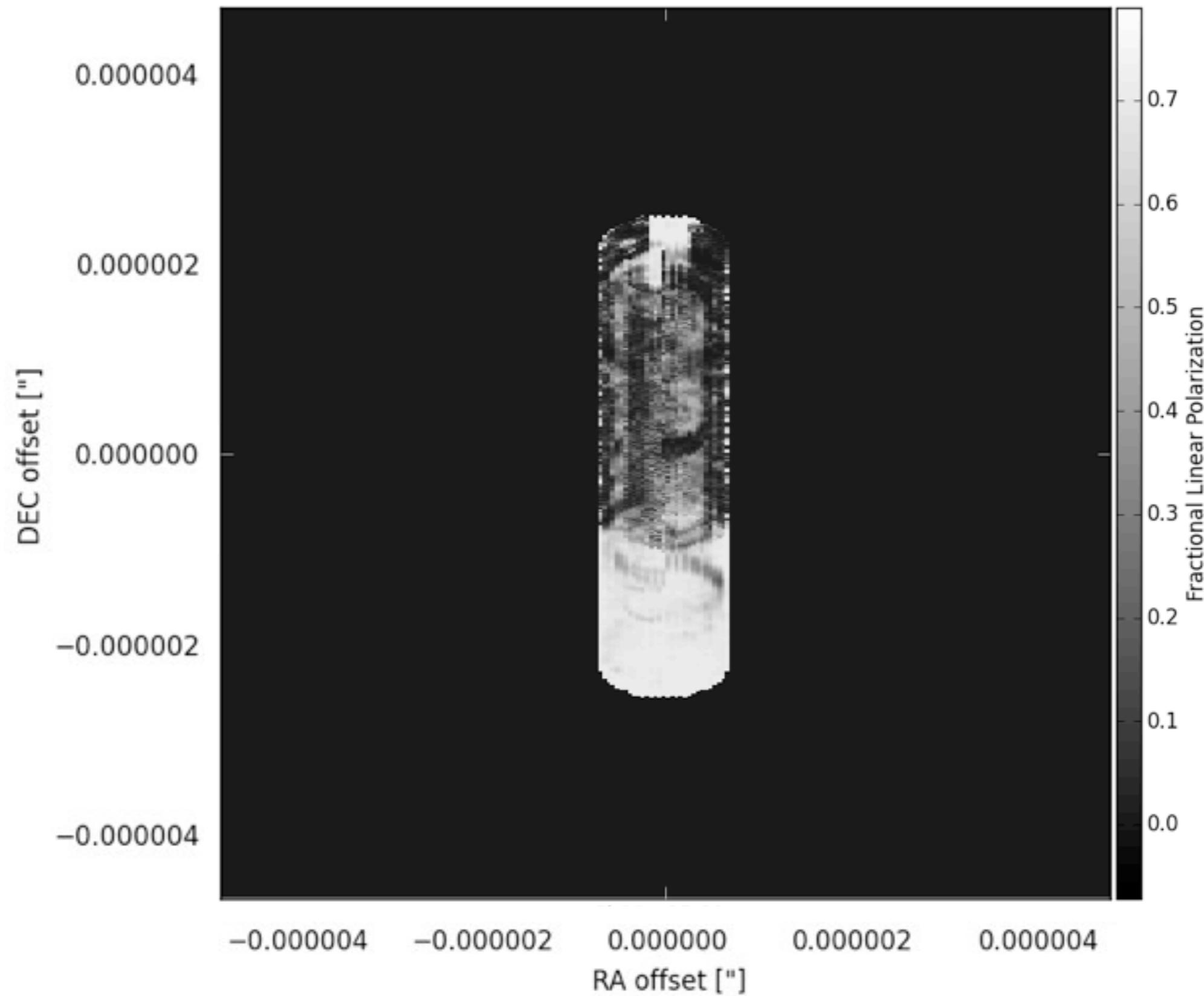
TEMZ Model (Disorder):



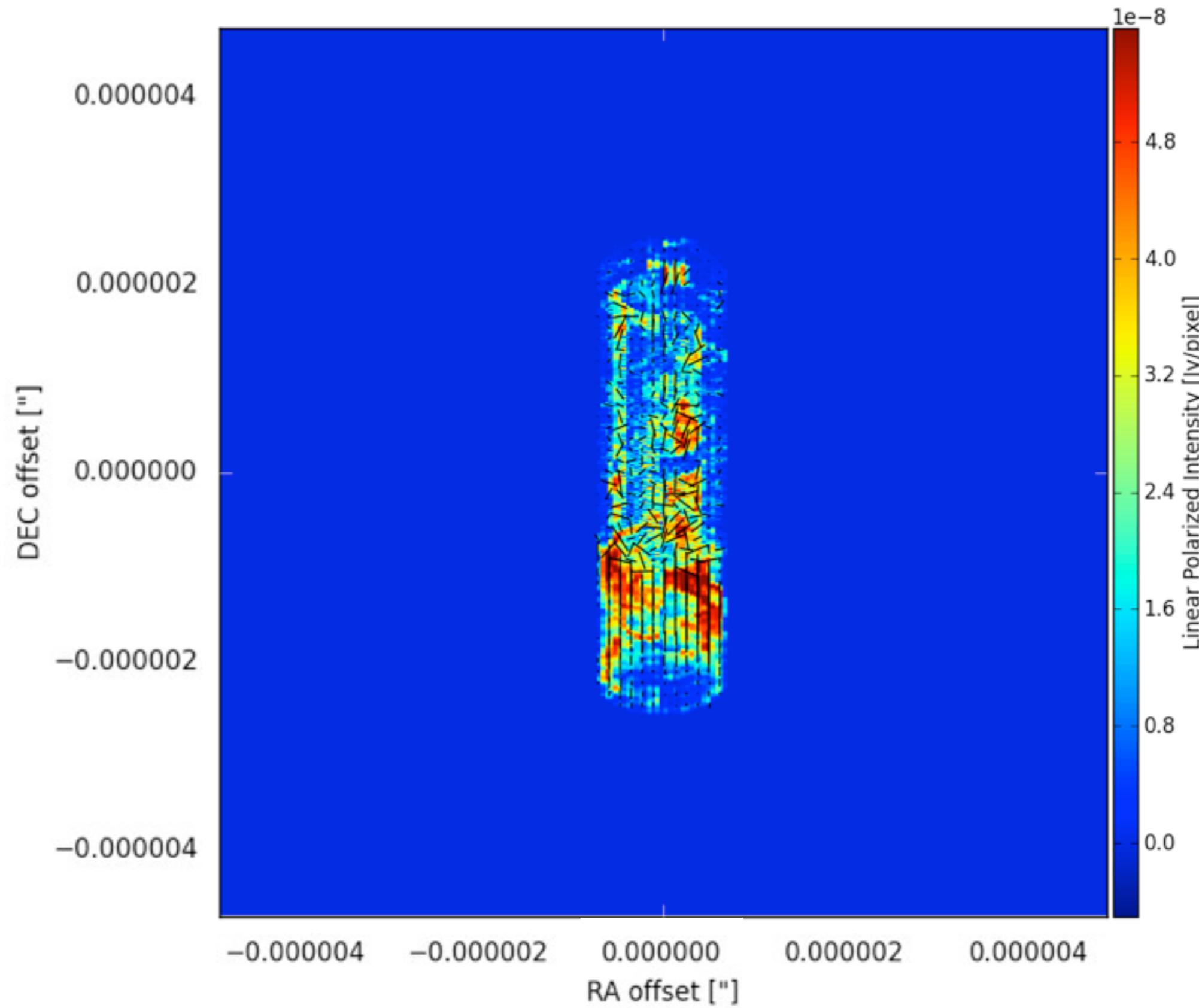
TEMZ Model (Disorder):



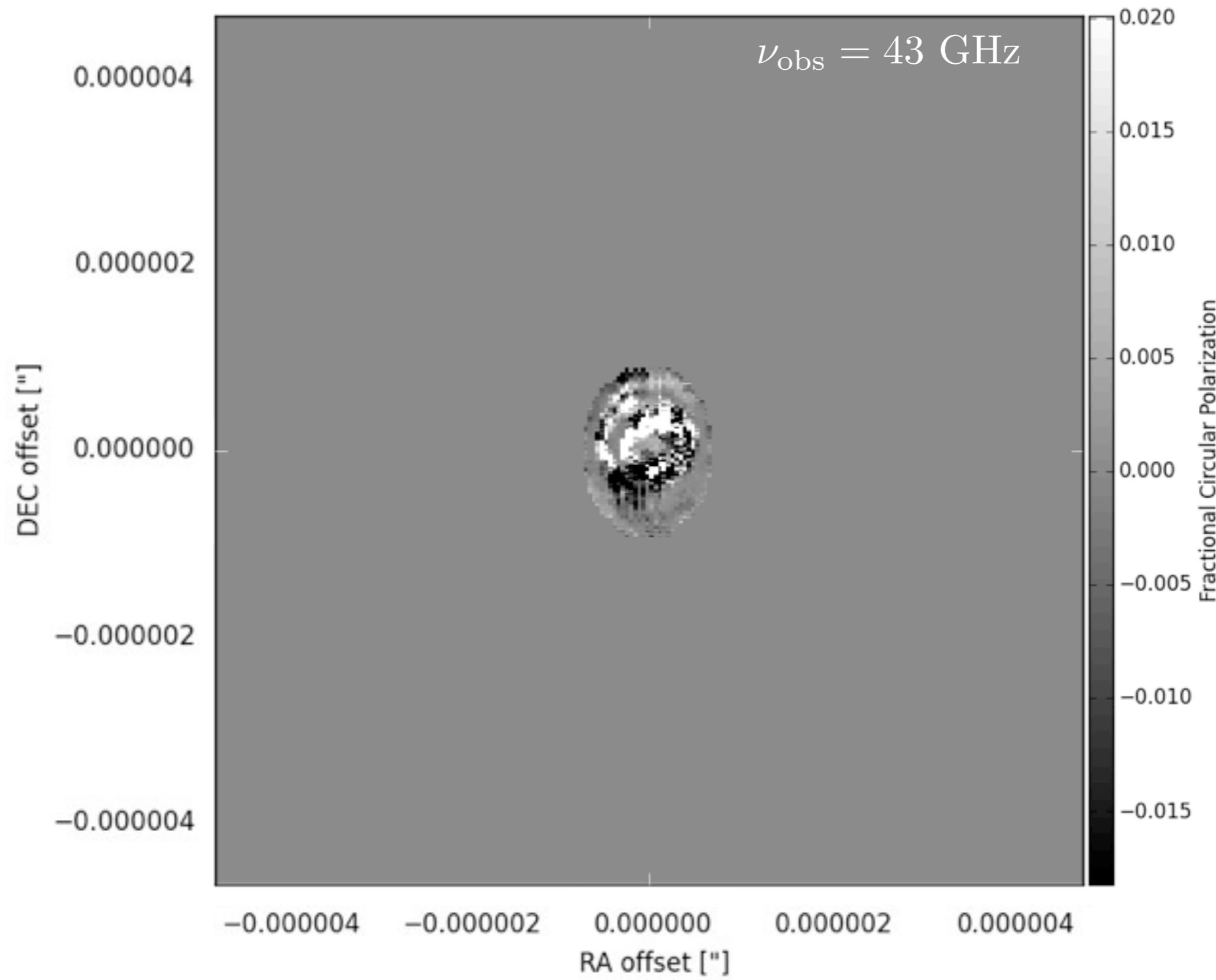
TEMZ Model (Disorder):



TEMZ Model (Disorder):



TEMZ Model (Disorder):



Can circular polarization be produced within a turbulent jet?



Questions?